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TESS Delivers Five New Hot Giant Planets Orbiting Bright Stars From The Full-Frame Images

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TESS Delivers Five New Hot Giant Planets Orbiting Bright Stars from the Full Frame Images

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ABSTRACT

We present the discovery and characterization of five hot and warm Jupiters — TOI-628 b (TIC 281408474; HD 288842), TOI-640 b (TIC 147977348), TOI-1333 b (TIC 395171208, BD+47 3521A), TOI-1478 b (TIC 409794137), and TOI-1601 b (TIC 139375960) — based on data from NASA's *Transiting Exoplanet Survey Satellite (TESS)*. The five planets were identified from the full frame images and were confirmed through a series of photometric and spectroscopic follow-up observations by the *TESS* Follow-up Observing Program (TFOP) Working Group. The planets are all Jovian size ($R_P = 1.01-1.77 R_J$) and have masses that range from 0.85 to 6.33 M_J. The host stars of these systems have F and G spectral types (5595 $\leq T_{eff} \leq 6460 K$) and are all relatively bright (9.5 < V < 10.8, 8.2 < K < 9.3) making them well-suited for future detailed characterization efforts. Three of the systems in our sample (TOI-640 b, TOI-1333 b, and TOI-1601 b) orbit subgiant host stars

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(log g <4.1). TOI-640 b is one of only three known hot Jupiters to have a highly inflated radius ($R_P > 1.7R_J$, possibly a result of its host star's evolution) and resides on an orbit with a period longer than 5 days. TOI-628 b is the most massive hot Jupiter discovered to date by *TESS* with a measured mass of $6.31^{+0.28}_{-0.30}$ M_J and a statistically significant, non-zero orbital eccentricity of $e = 0.074^{+0.021}_{-0.022}$. This planet would not have had enough time to circularize through tidal forces from our analysis, suggesting that it might be remnant eccentricity from its migration. The longest period planet in this sample, TOI-1478 b (P = 10.18 days), is a warm Jupiter in a circular orbit around a near-Solar analogue. NASA's *TESS* mission is continuing to increase the sample of well-characterized hot and warm Jupiters, complementing its primary mission goals.

1. INTRODUCTION

The discovery of hot Jupiters, combined with the assumption that gas giant planets must form at separations from their host star similar to our own giant planets, indicated that giant planets likely undergo large-scale migration from their formation locations. Various mechanisms have been proposed to emplace giant planets into very short-period orbits (Goldreich & Tremaine 1980; Lin & Papaloizou 1986; Lin et al. 1996; see Dawson & Johnson 2018 for a detailed review). However, it is not clear which of these mechanisms (if any) are dominant, govern this migration, or whether hot Jupiters can form in situ (Batygin et al. 2016), obviating the need for large-scale migration. One possibility is that giant planets migrate slowly and smoothly within the circumstellar gas-dust disk, resulting in well-aligned, nearly circular orbits (D'Angelo et al. 2003). It is also thought that planetary migration may be heavily influenced by gravitational interactions with other bodies within the system. These interactions result in highly eccentric and misaligned orbits (relative to the rotation axis of the star), and are typically referred to as a type of "High Eccentricity Migration" (HEM) or "Kozai-Lidov" (Kozai 1962; Lidov 1962; Rasio & Ford 1996; Wu & Murray 2003; Fabrycky & Tremaine 2007; Nagasawa & Ida 2011; Wu & Lithwick 2011; Naoz 2016). For short-period hot Jupiters, with periods less than about 5 days, the orbits will circularize in only a few billion years, erasing the evidence of HEM. Additionally, these interactions can cause misalignments in the planet's orbital plane (relative to the original disk plane), that can remain present for much longer, and such misalignments can be detected through Doppler spectroscopy, using observations of the Rossiter McLauglin effect (Rossiter 1924; McLaughlin 1924) or Doppler tomography (e.g., Miller et al. 2010; Johnson et al. 2014; Zhou et al. 2016). Longer-period hot (P \gtrsim 5 days) and warm (P>10 days) Jupiters experience smaller tidal forces, preserving their orbital eccentricities. It is likely that multiple mechanisms shape the short-period giant planet population, and studying these longer period hot Jupiters may give clues to their common evolutionary pathway.

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While some ground-based transit surveys that were dedicated to discovering hot Jupiters had near 24-hour coverage (Bakos et al. 2013), in general they struggled to discover planets with periods $\gtrsim 5$ days due to the poor duty cycle from weather and only being able to observe at night (Gaudi et al. 2005). Additionally, many of the first hot Jupiters discovered were assumed to reside in circular orbits when analyzing the observations, an assumption that may confuse current efforts to understand migration. Fortunately, NASA's Transiting Exoplanet Survey Satellite (TESS) mission was launched in April of 2018, and completed its primary 2-year long mission in July of 2020 (Ricker et al. 2015). TESS was awarded a 27-month first extended mission in which it will not only re-observe some areas covered in the primary mission but also observe most of the ecliptic plane, nearly completing coverage of the entire sky. TESS has a minimum observing baseline of ~ 27 days, and from recent occurrence rate studies, TESS planet searches will be mostly complete for hot Jupiters with periods ≤ 10 days (Zhou et al. 2019). Therefore, TESS provides a great resource for the discovery and confirmation of new longer-period hot and warm Jupiters (5 < P < 15 days) where eccentricities from migration would not be completely erased by tidal forces. TESS has already discovered a number of statistically significant hot Jupiters with highly eccentric (e > 0.2) orbits like HD 2685 b (Jones et al. 2019), TOI-172 b (Rodriguez et al. 2019), TOI-150 b (Kossakowski et al. 2019), TIC 237913194 b (Schlecker et al. 2020), and TOI-559 b (Ikwut-Ukwa et al. 2021). Additionally, TESS recently confirmed that the hot Jupiter HD 118203 b, an RV-identified planet with high eccentricity discovered with the radial velocity method (da Silva et al. 2006), transits its host star (Pepper et al. 2019).

TESS will also provide the ability to study hot Jupiter reinflation since its high photometric precision will allow it to discover giant planets around larger, more evolved host stars. As a star evolves off the main sequence, the stellar irradiation received by warm Jupiters is similar to that of a hot Jupiter. Therefore, discovering gas giant planets orbiting evolved stars at longer periods (10s of days) can test whether this increased irradiation causes the same inflation seen for short-period hot Jupiters (Lopez & Fortney 2016). Most warm Jupiters orbiting main sequence stars show little to no inflation (Demory & Seager 2011), suggesting that this energy must be transferred deep into the planet's interior (Liu

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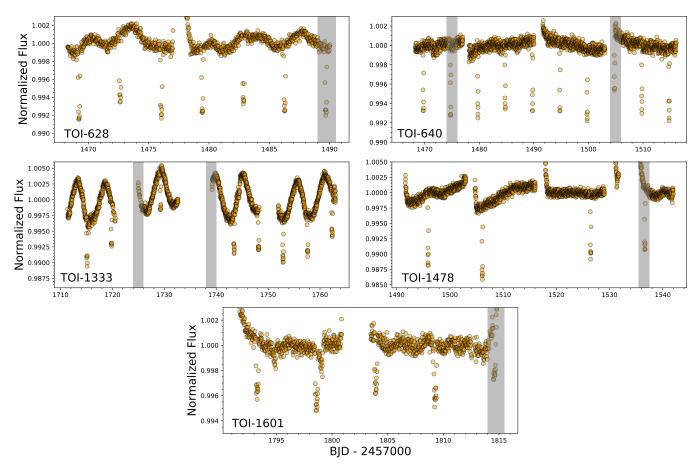


Figure 1. The raw *TESS* QLP 30-minute light curves for (top-left) TOI-628, TOI-640 (top-right), TOI-1333 (middle-left), TOI-1478 (middle-right), and TOI-1601 (bottom). Transits highlighted in gray were excluded from the global fit since they were flagged as bad quality by the QLP pipeline (Huang et al. 2020a).

et al. 2008; Spiegel & Burrows 2013) and as the star evolve, these warm Jupiters may re-inflate from the increased irradiation. Recent discoveries of hot Jupiters orbiting evolved stars are suggestive of reinflation (Almenara et al. 2015; Grunblatt et al. 2016; Hartman & Bakos 2016; Stevens et al. 2017) and *TESS* has already found a few hot and warm Jupiters orbiting evolved stars (Brahm et al. 2019; Nielsen et al. 2019; Huber et al. 2019; Rodriguez et al. 2019; Wang et al. 2019; Sha et al. 2020).

In this paper we present the discovery of TOI-628 b, TOI-640 b, TOI-1333 b, TOI-1478 b, and TOI-1601 b, five new hot and warm Jupiters from NASA's *TESS* mission. All five planets were discovered from an analysis of the 30-minute cadence Full Frame Images (FFIs), and first identified as *TESS* Objects of Interest (TOIs) by the *TESS* Science Office¹. These five new systems increase the known sample of wellcharacterized hot Jupiters, particularly those with longer orbital periods (>5 days). In §2 we present our time-series photometric and spectroscopic observations obtained by the TESS Follow-up Observing Program (TFOP) Working Group (WG) and describe the high spatial resolution imaging of all five targets, specifically on TOI-1333 and its two nearby companions. Our methodology for our global modeling using EXOFASTv2 (Eastman et al. 2019) is summarized in §3. We place these systems in context with the known population of hot and warm Jupiters, and discuss the impact of *TESS* in the discovery of giant planets in §7 and our conclusion are given in §5.

2. OBSERVATIONS AND ARCHIVAL DATA

We used a series photometric and spectroscopic observations to confirm and characterize the new systems, including high-spatial resolution imaging to rule out false positives, confirm them as bona-fide planets, and measure key parameters such as orbital eccentricity and the planet's density.

2.1. TESS Photometry

The initial detection of the new planets came from data collected by the *TESS* mission. *TESS* images the sky with a $24^{\circ}x96^{\circ}$ field of view and observes the same stars for about a month before moving on to observe a different region. Dur-

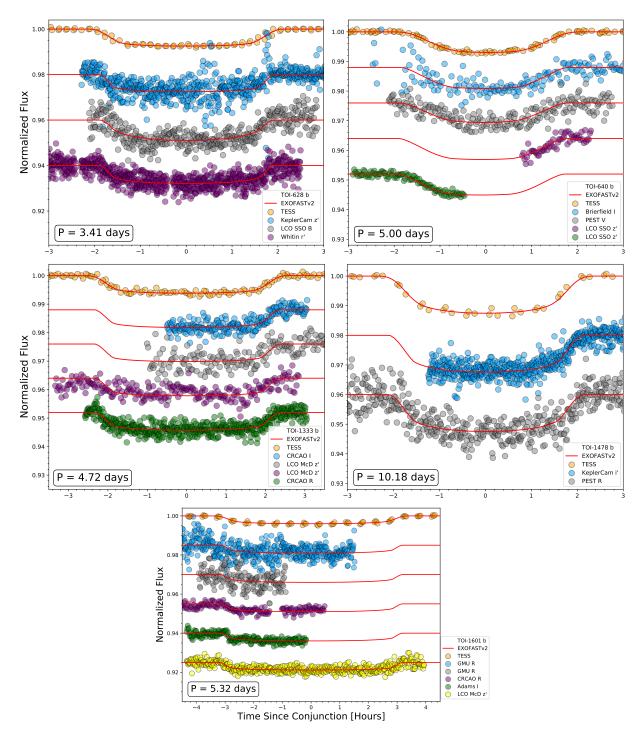


Figure 2. The *TESS* (orange) and TFOP SG1 follow-up transits of TOI-628 b (top-left), TOI-640 b (top-right), TOI-1333 b (middle-left), TOI-1478 b (middle-right), and TOI-1601 b (bottom). The EXOFASTv2 model for each transit observation is shown by the red solid line.

ing its primary mission, *TESS* saved and downloaded images of a smaller number of pre-selected stars every two minutes, while downloading co-added images of its entire field of view every 30 minutes. None of the planets described in this paper orbit stars that were pre-selected for two-minute cadence observations, so we use data from the 30-minute cadence FFIs.

After the data were transferred from the orbiting spacecraft back to Earth, the FFIs were calibrated using the TICA software (Fausnaugh et al. in prep) and light curves for a set of stars complete down to TESS-band magnitude = 13.5 were extracted with the MIT Quick Look Pipeline (QLP, Huang et al. 2020a). The QLP extracts light curves from the FFIs using a difference image analysis technique. After removing the effects of scattered earthshine and moonshine from the images, the QLP subtracts a high-quality reference image from each individual science frame and measures difference fluxes within sets of photometric apertures surrounding each star in the image. These difference fluxes are then converted to absolute brightness measurements by adding back the median flux expected from each star based on its TESSband magnitude. The QLP light curves have been used to discover dozens of planets (e.g., Huang et al. 2018; Rodriguez et al. 2019; Huang et al. 2020b) and a few thousand planet candidates (N. Guerrero et al. submitted). Additional information and a description of the QLP procedures are given by Huang et al. (2020a). The OLP lightcurves are shown in Figure 1. All five of these light curves were flattened using Keplerspline², a spline fitting routine to divide out the best-fit stellar variability (Vanderburg & Johnson 2014). The spacing of the spline break points for each system was determined by minimizing the Bayesian information criterion following the methodology from Shallue & Vanderburg (2018). After removing the stellar variability, we keep all data from one full transit duration before the transit until one full transit duration after the transit, and discard the remaining baseline data (which contains very little useful information but is computationally expensive to model). We use these lightcurve segments for the global fitting of each system (see §3).

Because the QLP uses aperture photometry on difference frames using a median combined reference image, it does not always measure absolute transit depths, only difference fluxes. Transit depths measured by QLP, especially in crowded fields, might be dependent on the accuracy of the *TESS*-band magnitudes from the *TESS* Input Catalog (Stassun et al. 2018) to correct for contamination due to blending of nearby stars within the aperture. To check that the transit depths measured by QLP for these targets are reasonably accurate, we performed a spot check for TOI-1601 with light curves extracted using a more traditional simple aperture photometry method (Vanderburg et al. 2019). We did not deblend the photometry in this test because TOI-1601 is in a relatively sparse field, as shown from our high resolution imaging (see §2.6.2). We ran an independent EXOFASTv2 fit for TOI-1601 b (swapping this lightcurve for the QLP one) following the strategy discussed in §3, and the results were consistent to within $<1\sigma$ uncertainties. Additionally, for each system we used the follow-up ground-based photometry within the global analysis to also constrain the depth, providing independent constraints that can be used to confirm the QLP depths (since the TFOP photometry is at a higher angular resolution than TESS). Within our global fit, we checked on any unknown contamination by fitting for a dilution term on the TESS bandpass to account for any difference compared to the SG1 photometry. In all cases other than TOI-1478, the fitted dilution was consistent with zero and well within our Gaussian 10% prior around zero, showing clear consistency between TESS and the TFOP seeing-limited photometry. TOI-1478 showed a significant required dilution on the order of 12%. To properly account for this, we removed the TESS dilution prior (see §3 for details), allowing it to be a free parameter, to properly correct for this within the fit.

We searched the non-flattened QLP light curves for rotation-based modulations using the VARTOOLS Lomb-Scargle function (Hartman & Bakos 2016). Specifically, we searched from 0.1 to 30 days and detect a clear strong periodicity at 5.296 days for TOI-1333, significantly different from the orbital period of its planetary companion ($P_b = 4.72$ days). This same periodicity is observed in our ground-based photometry, ruling out any systematics in the *TESS* observations. We see some tentative evidence of a periodicity at 10-11 days for TOI-628, however it was only observed in one *TESS* sector.

2.2. WASP Photometry

The WASP transit search consisted of two, wide-field arrays of eight cameras, with SuperWASP on La Palma covering the northern sky and WASP-South in South Africa covering the south (Pollacco et al. 2006). Each camera used a 200-mm, f/1.8 lens with a broadband filter spanning 400–700 nm, backed by 2048x2048 CCDs giving a plate scale of 13.7" pixel⁻¹. Observations then rastered available fields with a typical 15-min cadence.

We searched the WASP data for any rotational modulations using the methods from Maxted et al. (2011). TOI-640 was observed for spans of 150 nights in each of four years. The data from 2008, 2009 and 2010 show no significant modulation. The data from 2007, however, show significant power at a period of 62 ± 5 days, with an amplitude of 3 mmag and an estimated false-alarm probability below 1%. Since this is seen in only one season, and given that the data span only 2.5 cycles, we do not regard this detection as fully reliable.

² https://github.com/avanderburg/keplerspline

Table 1. Literature and Measured Properties

| Other identifiers | | | | | | | |
|------------------------------------|--|----------------------|----------------------|---------------------|----------------------|----------------------|------------------------|
| | | TOI-628 | TOI-640 | TOI-1333 | TOI-1478 | TOI-1601 | |
| | | TIC 281408474 | TIC 147977348 | TIC 395171208 | TIC 409794137 | TIC 139375960 | |
| | | HD 288842 | _ | BD+47 3521A | _ | — | |
| | TYCHO-2 | TYC 0146-01523-1 | TYC 7099-00846-1 | TYC 3595-01186-1 | TYC 5440-01407-1 | TYC 2836-00689-1 | |
| | 2MASS | J06370314+0146031 | J06385630-3638462 | J21400351+4824243 | J08254410-1333356 | J02332674+4100483 | |
| | TESS Sector | 6 | 6,7 | 15,16 | 7,8 | 18 | |
| Parameter | Description | Value | Value | Value | Value | Value | Source |
| α_{J2000} ‡ | Right Ascension (RA) | 06:37:03.13607 | 06:38:56.30742 | 21:40:03.50398 | 08:25:44.10708 | 02:33:26.74683 | 1 |
| δ_{J2000} \ddagger | Declination (Dec) | 01:46:03.19552 | -36:38:46.14425 | 48:24:24.52541 | -13:33:35.42756 | +41:00:48.36893 | 1 |
| B _T | Tycho B_T mag | 10.782 ± 0.051 | 11.178 ± 0.05 | 9.933±0.024 | 11.618 ± 0.082 | 11.521 ± 0.067 | 2 |
| $\mathbf{V}_T \dots \dots$ | Tycho V_T mag | $10.176 {\pm} 0.041$ | $10.574 {\pm} 0.043$ | 9.487 ± 0.021 | $10.805 {\pm} 0.067$ | $10.710 {\pm} 0.051$ | 2 |
| | Gaia <i>G</i> mag | $10.0579 {\pm} 0.02$ | $10.4006 {\pm} 0.02$ | $9.35 {\pm} 0.02$ | $10.66 {\pm} 0.02$ | $10.53 {\pm} 0.02$ | 1 |
| $B_P\ldots\ldots\ldots$ | Gaia B _P mag | $10.38 {\pm} 0.02$ | $10.68 {\pm} 0.02$ | $9.59 {\pm} 0.02$ | 11.03 ± 0.02 | $10.86 {\pm} 0.02$ | 1 |
| $R_P\ldots\ldots\ldots$ | Gaia R _P mag | $9.61 {\pm} 0.02$ | $9.99 {\pm} 0.02$ | $8.99 {\pm} 0.02$ | $10.16 {\pm} 0.02$ | $10.06 {\pm} 0.02$ | 1 |
| Τ | TESS mag | 9.6565 ± 0.0066 | 10.0367 ± 0.006 | 9.03527±0.0061 | 10.2042 ± 0.006 | 10.1035 ± 0.0065 | 3 |
| J | 2MASS J mag | 9.170±0.041 | 9.519±0.024 | 8.485±0.027 | 9.590±0.023 | 9.505±0.022 | 4 |
| Н | 2MASS H mag | $8.895 {\pm} 0.057$ | $9.327 {\pm} 0.026$ | $8.397 {\pm} 0.043$ | $9.255 {\pm} 0.026$ | $9.266 {\pm} 0.021$ | 4 |
| K _S | 2MASS K _S mag | 8.811 ± 0.02 | 9.243±0.023 | 8.272 ± 0.024 | 9.201 ± 0.021 | 9.19±0.02 | 4 |
| WISE1 | <i>WISE1</i> mag | 8.76±0.03 | 9.213±0.03 | 7.706±0.013 | 9.15±0.03 | 9.16±0.03 | 5 |
| WISE2 | <i>WISE2</i> mag | $8.79 {\pm} 0.03$ | $9.240 {\pm} 0.03$ | $7.594 {\pm} 0.012$ | $9.19 {\pm} 0.03$ | $9.21 {\pm} 0.03$ | 5 |
| WISE3 | <i>WISE3</i> mag | $8.59 {\pm} 0.03$ | $9.239 {\pm} 0.03$ | $8.035 {\pm} 0.019$ | $9.18 {\pm} 0.03$ | $9.18 {\pm} 0.034$ | 5 |
| WISE4 | <i>WISE4</i> mag | 8.23±0.19 | — | 8.043 ± 0.163 | _ | 8.80±0.33 | 5 |
| μ_{lpha} | Gaia DR2 proper motion in RA (mas yr ⁻¹) | -1.437±0.117 | -3.872±0.040 | -9.810±0.050 | -8.277±0.058 | 21.708±0.095 | 1 |
| μ_{δ} | Gaia DR2 proper motion in DEC (mas yr ⁻¹) | -1.082±0.087 | 4.927±0.045 | -10.501±0.046 | 67.943±0.046 | -0.874±0.090 | 1 |
| <i>v</i> sin <i>i</i> [*] | Rotational velocity (km s ⁻¹) | 6.9±0.5 | 6.1 ± 0.5 | 14.2±0.5 | 4.3±0.5 | 6.4±0.5 | §2.5.2& §2.5.2 & 2.5.3 |
| v _{mac} | macroturbulent broadening (km s ⁻¹) | $5.4 {\pm} 0.7$ | 6.32 ± 1.37 | 7.4 ± 1.8 | $4.9 {\pm} 0.5$ | 6.3 ± 0.6 | §2.5.2& §2.5.2 & 2.5.3 |
| π^{\dagger} | Gaia Parallax (mas) | $5.601 {\pm} 0.103$ | $2.925 {\pm} 0.033$ | $4.989 {\pm} 0.038$ | $6.542 {\pm} 0.047$ | $2.974 {\pm} 0.080$ | 1 |
| <i>P</i> _{Rot} | Rotation Period (days) | | | 5.3±0.159 | | | §2.1, 2.2, & 2.3 |

NOTES: The uncertainties of the photometry have a systematic error floor applied.

‡ RA and Dec are in epoch J2000. The coordinates come from Vizier where the Gaia RA and Dec have been precessed and corrected to J2000 from epoch J2015.5.

† Values have been corrected for the -0.30 μas offset as reported by Lindegren et al. (2018).

* U is in the direction of the Galactic center.

See D in the appendix of Collins et al. (2017) for a description of each detrending parameter. References are: ¹Gaia Collaboration et al. (2018), ²Høg et al. (2000), ³Stassun et al. (2018), ⁴Cutri et al. (2003), ⁵Zacharias et al. (2017)

For TOI-1333, the WASP data span 130 days in 2007, and show a clear modulation with a period of 15.9 ± 0.3 days, an amplitude of 19 mmag, and a false-alarm probability below 1%. A similar periodicity is also seen in the *TESS* data and in the Kilodegree Extremely Little Telescope (KELT) data (see below) but at a third of this period at 5.3 days. No significant periodicity was detected for TOI-1478 or TOI-1601.

Given knowledge of the *TESS* detections, transits of three of the systems described here can be found readily in the WASP data. TOI-640 was observed between 2006 and 2012, accumulating 23 000 data points. The WASP search algorithm (Collier Cameron et al. 2007) finds the transit with a period of 5.003773 ± 0.000041 and a mid-transit

epoch (T_C) of 2454822.00318 \pm 0.00411 HJD_{TDB}. This detection had been overlooked by WASP vetters owing to the near-integer day period (5.00 days), since the dominant red noise in WASP data is at multiples of a day. TOI-1478 was observed between 2009 and 2012, accumulating 9000 data points, less than usual for WASP since the field is near the crowded Galactic plane. It had not been flagged as a WASP candidate, however the search algorithm finds the transit with a period of 10.18051 \pm 0.00017 days and a T_C of 2455696.36710 \pm 0.00492 HJD_{TDB}. TOI-1601 was observed over 2006 and 2007, accumulating 10,400 data points. The search algorithm finds the transit and gets a period of 5.33197 \pm 0.00010 and a T_C of

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Table 2. Photometric follow-up observations of these systems used in the global fits and the detrending parameters.

| | | | • • • • | | 5011 | DI 10 1 | P () | |
|------------|----------------|-------------------|----------|--------|----------------------------------|------------------------|-------------|------------------------------|
| Target | Observatory | Date (UT) | size (m) | Filter | FOV | Pixel Scale | Exp (s) | Additive Detrending |
| TOI-628 b | FLWO/KeplerCam | 2019 December 06 | 1.2 | z' | $23.1' \times 23.1'$ | 0.672'' | 9 | None |
| TOI-628 b | LCO TFN | 2019 December 16 | 0.4 | z' | 19' 	imes 29' | 0.57'' | 18 | airmass |
| TOI-628 b | Whitin | 2020 February 08 | 0.7 | r' | $24' \times 24'$ | 0.67'' | 10 | airmass |
| TOI-640 b | Brierfield | 2019 December 27 | 0.36 | Ι | $49.4^\prime \times 49.4^\prime$ | 1.47'' | 90 | X(FITS), Y(FITS) |
| TOI-640 b | PEST | 2020 March 01 | 0.3048 | R | $31^\prime 	imes 21^\prime$ | 1.2'' | 60 | none |
| TOI-640 b | LCO SSO | 2020 August 23 | 1.0 | z' | $27' \times 27'$ | 0.39" | 60 | Y(FITS) |
| TOI-640 b | LCO SSO | 2020 November 06 | 1.0 | z' | $27' \times 27'$ | 0.39″ | 60 | airmass |
| TOI-1333 b | CRCAO | 2020 July 29 | 0.6096 | Ι | $26.8^\prime \times 26.8^\prime$ | 0.39" | 45 | airmass |
| TOI-1333 b | LCO McDonald | 2020 July 29 | 0.4 | z' | $19^\prime 	imes 29^\prime$ | 0.57'' | 30 | total counts |
| TOI-1333 b | LCO McDonald | 2020 August 12 | 0.4 | z' | 19' 	imes 29' | 0.57'' | 30 | Y(FITS) |
| TOI-1333 b | CRCAO | 2020 September 19 | 0.6096 | R | $26.8^\prime \times 26.8^\prime$ | 0.39" | 80 | airmass |
| TOI-1478 b | FLWO/KeplerCam | 2019 December 14 | 1.2 | i' | $23.1^\prime \times 23.1^\prime$ | $0.672^{\prime\prime}$ | 7 | none |
| TOI-1478 b | PEST | 2020 January 03 | 0.3048 | R | $31^\prime 	imes 21^\prime$ | 1.2'' | 30 | airmass |
| TOI-1601 b | GMU | 2020 August 30 | 0.8 | R | 23' 	imes 23' | 0.34″ | 30 | airmass, sky/pixels |
| TOI-1601 b | GMU | 2020 September 15 | 0.8 | R | 23' 	imes 23' | 0.34″ | 30 | airmass, sky/pixels, X(FITS) |
| TOI-1601 b | CRCAO | 2020 October 01 | 0.6096 | z' | $26.8^\prime \times 26.8^\prime$ | 0.39" | 90 | airmass |
| TOI-1601 b | Adams | 2020 October 17 | 0.61 | Ι | $26^\prime 	imes 26^\prime$ | 0.38″ | 60 | airmass, total counts |
| TOI-1601 b | LCO McDonald | 2020 October 17 | 1.0 | z' | $27' \times 27'$ | 0.39" | 60 | airmass, sky/pixels |

NOTES: All the follow-up photometry presented in this paper is available in machine-readable form in the online journal.

2454186.65253 \pm 0.01283 HJD_{TDB}. We use these T_C values as priors for the EXOFASTv2 global fits of TOI-640 b, TOI-1478 b, and TOI-1601 b. We see a ~40% reduction in uncertainty on the period of the planet when including the WASP T_C prior.

2.3. KELT Photometry

To complement the TESS photometry, we analyzed observations of these five TOIs from the KELT survey³ (Pepper et al. 2007, 2012, 2018). For a full description of the KELT observing strategy and reduction process, see Siverd et al. (2012); Kuhn et al. (2016). KELT has two fully robotic telescopes, each of which uses a Mamiya 645 80mm f/1.9, each of which uses a Mamiva 645 80mm f/1.9 lens with 42mm aperture and Apogee 4k×4k CCD on a Paramount ME mount. This setup provides to a $26^{\circ} \times 26^{\circ}$ field of view with a 23'' pixel scale. The two telescopes are located in at Winer Observatory in Sonoita, AZ and at the South African Astronomical Observatory (SAAO) in Sutherland, South Africa. KELT covers \sim 85% of the entire sky and the observing strategy results in a 20-30 minute cadence. Some of the KELT lightcurves are publicly available through the NASA Exoplanet Archive. KELT observations were available for TOI-628, TOI-1333, TOI-1478, and TOI-1601. KELT-South observed TOI-628 2828 times and TOI-1478 4632 times from

2010 to 2015. KELT-North observed TOI-1333 from 2012 to 2014 and TOI-1601 from 2006 to 2014, acquiring 2580 and 8520 observations, respectively. Since KELT has been observing since 2006 in some cases, the observations significantly extend the baseline of the photometry and can provide a strong constraint on the ephemeris of each system. Following the strategy described in Siverd et al. (2012) and Kuhn et al. (2016), we also search the KELT lightcurves for transits of each planet. Unfortunately, no significant signs of the known planetary transits were found, likely due to the poor duty cycle for longer orbital periods.

Following the approach of Stassun et al. (1999) and Oelkers et al. (2018), we executed a search for periodic signals most likely to come from the rotation period of the star. For these stars, we post-processed the light curve data using the Trend-Filtering Algorithm (Kovács et al. 2005) to remove common systematics. We then searched for candidate rotation signals using a modified version of the Lomb-Scargle period finder algorithm (Lomb 1976; Scargle 1982). We searched for periods between a minimum period of 0.5 days and a maximum period of 50 days using 2000 frequency steps⁴. We masked periods between 0.5 and 0.505 days and 0.97–1.04 days to avoid the most common detector aliases associated with KELT's observational cadence and its inter-

⁴ The total number of frequency steps may vary slightly depending on the number of data points in any given light curve.

action with the periods for the solar and sidereal day. For each star, we selected the highest statistically significant peak of the power spectrum (hereafter, γ) as the candidate period.

We then executed a boot-strap analysis, using 1000 Monte-Carlo iterations, where the dates of the observations were not changed, but the magnitude values of the light curve were randomized, following the work of Henderson & Stassun (2012). We recalculated the Lomb-Scargle power spectrum for each iteration, and recorded the maximum peak power (hereafter, γ_{sim}) of all iterations. If the highest power spectrum peak was larger than the maximum simulated peak ($\gamma > \gamma_{sim}$) after 1000 iterations, we considered the periodic signal to be a candidate rotation period.

We find TOI-1333 to have a strongly significant ($\gamma > 50$) candidate rotation period at 5.3 days, which is consistent with *TESS* (see §2.1), and TOI-1601 and TOI-1478 to have weakly significant ($\gamma > 10$) candidate rotation periods of 9.3 and 16.6 days, respectively. However, we do not see these periodicities in the *TESS* photometry, and they are likely aliases of the KELT observing strategy. KELT did not obtain observations of TOI-640.

2.4. Ground-based Photometry from the TESS Follow-up Observing Program Working Group

To refine the ephemerides and transit parameters for each system while ruling out false positive scenarios, we obtained photometric follow-up observations on all five systems from the TESS Follow-up Observing Program Working Group⁵, sub-group 1 (SG1) for seeing-limited photometry. Specifically, the follow-up comes from the Las Cumbres Observatory (LCO) telescope network (Brown et al. 2013), Whitin Observatory at Wellesley College, KeplerCam on the 1.2m telescope at Fred Lawrence Whipple Observatory (FLWO), Brieffield Observatory, PEST Observatory, C. R. Chambliss Astronomical Observatory (CRCAO) at Kutztown University, Adams Observatory at Austin College, and Suto Observatory. To schedule the photometric transit follow-up observations, we used the TAPIR software package (Jensen 2013). The data reduction and aperture photometry extraction was performed using AstroImageJ (Collins et al. 2017) for all follow-up data other than the PEST observations which were done using a custom software package PEST Pipeline⁶. The TFOP follow-up photometry for these five systems are shown in Figure 2, and a list of each telescope's information and details on each follow-up transit can be seen in Table 2. The follow-up transits presented in Figure 2 are available as machine-readable tables with this paper.

2.5. Spectroscopy

Table 3. The first RV point from each spectrograph for all five systems. The full table of RVs for each system is available in machine-readable form in the online journal.

| BJD _{TDB} | RV (m s ⁻¹) | $\sigma_{RV}^{\dagger} (\mathrm{m}\mathrm{s}^{-1})$ | Target | Instrument |
|--------------------|-------------------------|---|----------|------------|
| 2458604.64716 | 193.0 | 43.6 | TOI-628 | TRES |
| 2458742.25270 | 19775.2 | 17.5 | TOI-628 | MINERVA3 |
| 2458742.25270 | 19700.8 | 14.3 | TOI-628 | MINERVA4 |
| 2458742.25270 | 19795.7 | 41.0 | TOI-628 | MINERVA6 |
| 2458748.89482 | 39168.0 | 29.4 | TOI-640 | CHIRON |
| 2458777.69583 | -50.6 | 48.0 | TOI-1333 | TRES |
| 2458852.80172 | 19374.4 | 21.9 | TOI-1478 | CHIRON |
| 2458875.60994 | 20738.5 | 22.0 | TOI-1478 | CORALIE |
| 2458910.73537 | 20910.3 | 7.5 | TOI-1478 | FEROS |
| 2458829.86525 | 67.4 | 27.6 | TOI-1478 | TRES |
| 2459184.83629 | 19502.0 | 24.0 | TOI-1478 | CHIRON2 |
| 2458847.77575 | -184.1 | 32.0 | TOI-1601 | TRES |
| | | | | |

NOTES: † The internal RV error for the observation shown.

To rule out false positive scenarios and measure the mass and orbital eccentricity of each system, we obtained time series spectroscopy coordinated through the TFOP WGs. A sample of one radial velocity (RV) point per target per instrument is shown in Table 3, with the full table available in machine-readable form in the online journal. The RVs and best-fit models from our EXOFASTv2 analysis are shown in Figure 3 (see §3). Following the methodology in Zhou et al. (2018), we measure the $v \sin I_*$ and macroturbulent broadening for all five systems from TRES except TOI-640 where the CHIRON observations were used (see Table 1).

2.5.1. TCES Spectroscopy

Reconnaissance spectroscopic observations of TOI-628 were carried out with the Tautenburg Coudé Echelle Spectrograph (TCES) mounted at the 2-meter "Alfred Jensch" Telescope of the Thuringian State Observatory (TLS) in Tautenburg, Germany. The spectra cover the 470-740 nm wavelength range and have a resolution R=57000. A 40 min exposure was taken at BJD=2458777.6053 (orbital phase $\phi \sim 0.5$), and a 3×20 min exposure at BJD=2458855.3720 ($\phi \sim 0.0$). We measured for the two single-lined spectra a $\Delta RV \leq 1$ km s⁻¹, ruling out an eclipsing binary as the cause of the event detected by TESS. These velocities were not included in the global fit for TOI-628.

2.5.2. TRES Spectroscopy

To confirm targets from *TESS* in the Northern hemisphere, we observed TOI-628, TOI-1333, TOI-1478, and TOI-1601 with the Tillinghast Reflector Echelle Spectrograph (TRES;

⁵ https://tess.mit.edu/followup/

⁶ http://pestobservatory.com/the-pest-pipeline/

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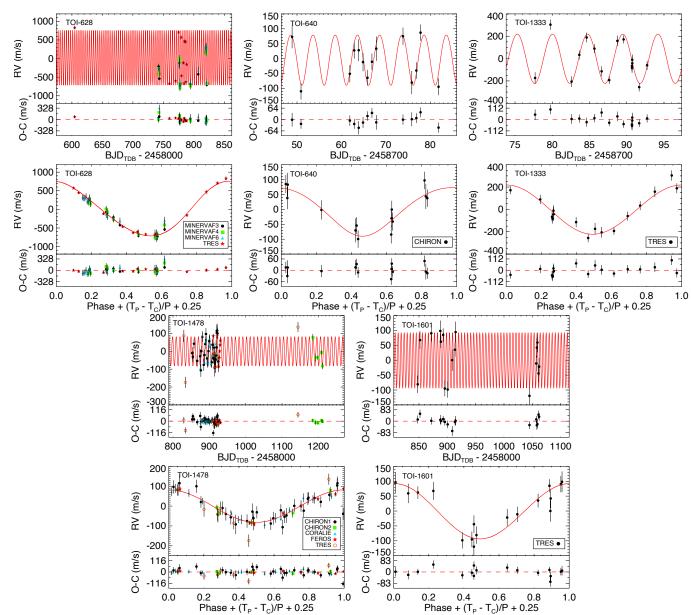


Figure 3. The RV observations of TOI-628 (top-left), TOI-640 (top-middle), TOI-1333 (top-right), TOI-1478 (bottom-left), and TOI-1601 (bottom-right). In each case, the top figure shows the RVs vs time and the bottom panel is phased to the best-fit ephemeris from our global fit. The EXOFASTv2 model is shown in red and the residuals to the best-fit are shown below each plot. We see no periodicity in the residuals from our fit.

Fűrész 2008)⁷ on the 1.5m Tillinghast Reflector. The telescope is located at the Fred L. Whipple Observatory (FLWO) on Mt. Hopkins, AZ, and the spectrograph has a resolving power of R=44,000. See Buchhave et al. (2010) and Quinn et al. (2012) for a detailed description on reduction and RV extraction pipeline. The only difference in our analysis is that we created the template spectra for the RV extraction by aligning and median-combining all of the out-of-transit spectra. We removed cosmic rays and cross-correlated the median combined spectra against all the observed spectra. Bisector spans for the TRES spectra were calculated following the technique described in Torres et al. (2007). There was no correlation between the bisector spans and the RVs. We also used the TRES spectra to provide constraints on the $T_{\rm eff}$ and [Fe/H] for our global analysis. We analyzed the TRES spectra with the Stellar Parameter Classification (SPC) package (Buchhave et al. 2012) to determine the metallicity and rotational velocity for all five host stars (see Table 1 and 4). We also used SPC to determine a constraint on the $T_{\rm eff}$ of 6250±100 K for TOI-1333, which is used in §2.6.4 to con-

⁷ http://www.sao.arizona.edu/html/FLWO/60/TRES/GABORthesis.pdf

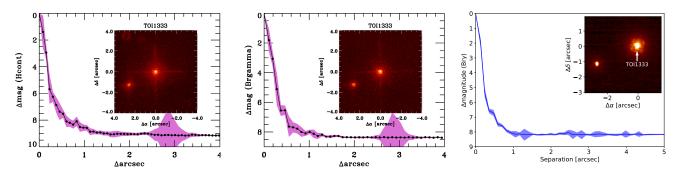


Figure 4. The Palomar PHARO (left) H-band and (middle) Br γ -band 4 σ contrast curve for TOI-1333 with the AO image embedded in the plot. The (right) Gemini NIRI Br γ -band AO 5 σ contrast curve for TOI-1333. The NIRI AO image is embedded in the plot. The second star in the image is TIC 2010985858, and we properly account for its blending in our fit (see §3). The colored swath represents the uncertainty on the 5 σ contrast curve (see §2.6.2).

strain the dilution from nearby companions and the radius of TOI-1333.

2.5.3. CHIRON Spectroscopy

We obtained a series of spectroscopic observations with the 1.5 m SMARTS / CHIRON facility (see Table 3 Tokovinin et al. 2013) for TOI-640 and TOI-1478 to measure the host star parameters, and constrain their masses and eccentricities. The 1.5 m SMARTS facility is located at Cerro Tololo Inter-American Observatory (CTIO), Chile. CHIRON is a high resolution echelle spectrograph fed via an image slicer through a single multi-mode fiber, with a spectral resolving power of $R \sim 80,000$ over the wavelength region from 410 to 870nm. For the case of TOI-1478, we treat the pre- and post-shutdown RVs as separate instruments within the fit (see §3).

To obtain the stellar atmospheric parameters, we matched the CHIRON spectra against an interpolated library of observed $\sim 10,000$ spectra classified by SPC (Buchhave et al. 2012). The metallicity from this analysis was used as a prior for global fit of TOI-640 (see Table 4). RVs were derived via the least-squares deconvolution (Donati et al. 1997; Zhou et al. 2020) of the spectra against non-rotating synthetic templates matching the spectral parameters of each host star, generated using the ATLAS9 model atmospheres (Kurucz 1992). The RVs were then measured by fitting a least-squares deconvolution line profile with a rotational broadening kernel as prescribed by Gray (2005). The velocities for each system are presented in Table 3.

2.5.4. FEROS Spectroscopy

Using the FEROS spectrograph (Kaufer et al. 1999) mounted on the 2.2m MPG telescope at La Silla observatory in Chile, we obtained 8 observations of TOI-1478. FEROS has R=48000, and the observations were between UT 2020 March 02 and UT 2020 March 22. A ThAr+Ne lamp was used to illuminate the fiber simultaneously to the observations to determine the instrumental offset. We reduced the spectra, derived the RVs, and produced the bisector spans

using the CERES suite for echelle pipelines (Brahm et al. 2017).

2.5.5. CORALIE Spectroscopy

TOI-1478 was observed with the CORALIE high resolution spectrograph (R=60,000) on the Swiss 1.2 m Euler telescope at La Silla Observatories, Chile (Queloz et al. 2001). A total of 14 spectra were obtained between UT 2020 January 26 and March 16, covering several orbits of TOI-1478 b. CORALIE is fed by a 2"science fiber and a secondary fiber with simultaneous Fabry-Perot for wavelength calibration. RVs were computed with the standard CORALIE data reduction pipeline via cross-correlation with a binary G2 mask. Activity indices, bisector spans (BIS) and other line profile diagnostics were extracted as well. We find no correlation between the RVs and BIS, nor activity indicators. Our exposure times varied between 1200 and 1800 seconds depending on site conditions and observing schedule.

2.5.6. MINERVA Australis Spectroscopy

MINERVA-Australis is an array of four PlaneWave CDK700 telescopes located in Queensland, Australia, fully dedicated to the precise RV follow-up of TESS candidates. The four telescopes can be simultaneously fiber-fed to a single KiwiSpec R4-100 high-resolution (R=80,000) spectrograph (Barnes et al. 2012; Addison et al. 2019, 2020). TOI-628 was monitored by MINERVA-Australis using 2 or 3 telescopes in the array (Minerva3, Minerva4, Minerva6) between UT 2019 Sep 15 and Dec 3. Each epoch consists of two 30minute exposures, and the resulting RVs are binned to a single point as shown by the example in Table 3. Telescopes 1 and 3 obtained seven RV epochs, while Telescope 4 obtained four epochs. RVs for the observations are derived for each telescope by cross-correlation, where the template being matched is the mean spectrum of each telescope. The instrumental variations are corrected by using the RVs computed from different Minerva telescopes as originating from independent instruments within our global model.

2.6. High Resolution Imaging

As part of our standard process for validating transiting exoplanets to assess possible contamination of bound or unbound companions on the derived planetary radii (Ciardi et al. 2015), we obtained high spatial resolution imaging observations of all five systems.

2.6.1. Speckle Imaging

We searched for close companions to TOI-628, TOI-640, and TOI-1478 with speckle imaging in the *I*-band on the 4.1-m Southern Astrophysical Research (SOAR) telescope (Tokovinin 2018). The speckle imaging was conducted using HRCam (field-of-view of 15") and had a 0.01575" pixel scale. TOI-628 was observed on UT 2019 November 11 with a sensitivity of Δ Mag = 7.2 at 1". Speckle observations of TOI-640 were taken on UT 2019 May 18 and had an estimated contrast of Δ mag = 6.6 at 1". Observations of TOI-1478 were taken on UT 2020 January 07 and had an estimated contrast of Δ mag = 6.8 at 1". See Ziegler et al. (2020) for a description on the general observing strategy for *TESS* targets. No nearby companion was observed for any of the three targets out to 3".

We also obtained two sets of high-resolution speckle images of TOI-1478. One was collected on UT 2020 January 14 using the Zorro instrument mounted on the 8-meter Gemini South telescope located on the summit of Cerro Pachon in Chile and the other on UT 2020 February 18 using the 'Alopeke instrument mounted on the 8-m Gemini-North telescope located on the summit of Mauna Kea in Hawaii. These twin instruments simultaneously observe in two bands, i.e., $\frac{\lambda}{\Delta\lambda}$ = 832/40 nm and 562/54 nm, obtaining diffraction limited images with inner working angles 0.026" and 0.017", respectively. Each observation consisted of 6 minutes of total integration time at each telescope taken as sets of 1000×0.06 second images. All the images were combined and subjected to Fourier analysis leading to the production of final data products including speckle reconstructed imagery (see Howell et al. 2011). Both speckle imaging results showed similar contrast limits and revealed that TOI-1478 is a single star to contrast limits of 5 to 9 magnitudes (out to 1.17'', ruling out most main sequence companions to TOI-1478 within the spatial limits of \sim 4 to 180 au (for d=153 pc, as determined from the Gaia DR2 parallax Gaia Collaboration et al. 2018).

2.6.2. Adaptive Optics Imaging

We observed TOI-628 on UT 2019 November 11 using the ShARCS adaptive optics system on the 3-m Shane Telescope at Lick Observatory. ShARCS has a field of view of $20'' \times 20''$ and a pixel scale of 0.033'' pixel⁻¹. We conducted our observations using a square 4-point dither pattern with a separation of 4'' between each dither position. Our observations were taken in natural guide star mode with high winds. We

obtained one sequence of observations in Ks-band and a second sequence in J-band for a total integration time of 510 s in Ks-band and 225 s in J-band. See Savel et al. (2020) for a detailed description of the observing strategy. Neither set of observations revealed any companions for TOI-628.

We observed TOI-1333 ($Br - \gamma$ and H - cont) and TOI-1478 with infrared high-resolution adaptive optics (AO) imaging at Palomar Observatory. The Palomar Observatory observations were made with the PHARO instrument (Hayward et al. 2001) behind the natural guide star AO system P3K (Dekany et al. 2013). The observations were made on 2019 Nov 10 UT in a standard 5-point quincunx dither pattern with steps of 5''. Each dither position was observed three times, offset in position from each other by 0.5'' for a total of 15 frames. The camera was in the narrow-angle mode with a full field of view of $\sim 25''$ and a pixel scale of approximately 0.025" per pixel. Observations were made in the narrow-band Br- γ filter ($\lambda_{\rho} = 2.1686; \Delta \lambda = 0.0326 \mu m$) for TOI-1333 and TOI-1478, and in the *H*-cont filter ($\lambda_{0} =$ 1.668; $\Delta \lambda = 0.018 \mu m$) for TOI-1333. The observations get down to a Δ Mag = 6.54 (Br- γ) and = 7.52 (*H*-*cont*) for TOI-1333 and a Δ Mag = 6.8 (Br- γ) for TOI-1478 (all at $\sim 0.5''$).

TOI-1333 was also observed using NIRI on Gemini-North (Hodapp et al. 2003) on UT 2019 November 14 in the Br- γ filter. NIRI has a 22" × 22" field-of-view with a 0.022" pixel scale. Our sequence consisted of nine images, each with exposure time 4.4s, and we dithered the telescope between each exposure. A sky background was constructed from the dithered frames, and subtracted from each science image. We also performed bad pixel removal and flatfielding, and then aligned and coadded frames. NIRI got down to a Δ Mag = 6.7 at 0.472".

We also observed TOI-1601 using the Near Infrared Camera 2 (NIRC2) adaptive optics (AO) set up on the W. M. Keck Observatory in the Br- γ filter and in the *J*-*cont* filter on UT 2020 September 09. The NIRC2 detector has a 9.971 mas pixel⁻¹ using a 1024×1024 CCD (field-of-view = 10" × 10", Service et al. 2016). Unfortunately, the lower left quadrant of the CCD is known to have higher than typical noise levels in comparison to the others. To avoid this issue, a 3-point dither pattern technique was used. The images were aligned and stacked after normal flat-field and sky background corrections. No nearby companions were seen down to a Δ mag = 6.680 (*J*-*cont*) and = 6.402 (Br- γ) for TOI-1601 at 0.5".

While the observing strategy differed, all of the AO data were processed and analyzed with a custom set of IDL tools. The science frames were flat-fielded and sky-subtracted. The flat fields were generated from a median combination of the dark subtracted flats taken on-sky. The flats were normalized such that the median value of the flats is unity. The sky frames were generated from the median average of the 15 dithered science frames; each science image was then

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Table 4. Median values and 68% confidence intervals for the global models

| Priors: | | TOI-628 b | TOI-640 b | TOI-1333 b§ | TOI-1478 b¶ | TOI- | 1601 b |
|-----------------------------|--|--|---|--|--|---|--|
| Gaussian | π Gaia Parallax (mas) | 5.601±0.103 | 2.925±0.033 | _ | 6.542±0.047 | 2.974 | ± 0.080 |
| Gaussian | [Fe/H] Metallicity (dex) | 0.24 ± 0.08 | 0.02 ± 0.1 | 0.11 ± 0.08 | 0.07 ± 0.08 | | ± 0.08 |
| Upper Limit | A_V V-band extinction (mag) | 2.977 | 0.292 | _ | 0.120 | | 4694 |
| Gaussian | R_{\star} Stellar Radius (R_{\odot}) | _ | _ | 1.963 ± 0.064 | _ | _ | _ |
| Gaussian | T _{eff} Stellar Effective Temperature (K) | _ | _ | 6250 ± 100 | _ | _ | _ |
| Gaussian | T_{C}^{**} Time of conjunction (HJD _{TDB}) | | 2454822.00318±0.00411 | 0250 ± 100 | 2455696.36710±0.00492 | 2454186.653 | 253±0.01283 |
| Taussian/ | D_T Dilution in TESS | - 0.00000 \pm 0.00344 | 0.00000±0.00360 | - 0.00000 \pm 0.03817 | 2455090.50710±0.00492 | | 0.00196 |
| Jaussian Gaussian*** | • | 0.00000 ± 0.00344 | 0.00000 ± 0.00360 | | — | 0.0主(| 0.00196 |
| | D _I Dilution in I | — | _ | 0.3609±0.0180 | — | - | _ |
| Gaussian*** Gaussian*** | D_R Dilution in R $D_{a'}$ Dilution in z' | _ | _ | 0.3499 ± 0.0175 0.0664 ± 0.0033 | — | - | _ |
| arameter | Units | Values | — | 0.0004_0.0055 | _ | | |
| arameter tellar Paramete | | values | | | | | |
| robability | | 100 % | 100 % | 100 % | 100% | 68.4 % | 31.6% |
| M* | Mass (M ₍₎) | 1.311 ^{+0.066} -0.075 | 1.536 ^{+0.069} -0.076 | 1.464 ^{+0.076} -0.079 | 0.947 ^{+0.059} -0.041 | 1.517 ^{+0.053} -0.049 | 1.340 ^{+0.042} -0.045 |
| | 0 | 1.511-0.075 | $^{1.550}$ -0.076 | 1.404 - 0.079 | 0.947-0.041 | | 1.340-0.045 |
| <i>R</i> * | Radius (<i>R</i> _☉) | $1.345^{+0.046}_{-0.040}$ | 2.082+0.064 -0.058 | $1.925^{+0.064}_{-0.063}$ | $1.048 \substack{+0.030 \\ -0.029}$ | $2.186^{+0.074}_{-0.063}$ | $2.185^{+0.086}_{-0.080}$ |
| <i>L</i> * | Luminosity (L_{\odot}) | 2.48 ^{+0.41} -0.28 | 6.83 ^{+0.41} -0.54 | 5.17 ^{+0.49} -0.45 | $0.971^{+0.036}_{-0.035}$ | 5.40 ^{+0.35} -0.33 | 5.25 ^{+0.35} -0.33 |
| F_{Bol} | Bolometric Flux $\times 10^{-9}$ (cgs) | 2.52 ^{+0.39} -0.27 | $1.87^{+0.10}_{-0.14}$ | _ | 1.32 ± 0.45 | $1.505^{+0.52}_{-0.61}$ | 1.485 ^{+0.59} -0.61 |
| ρ_* | Density (cgs) | $0.762^{+0.067}_{-0.078}$ | $0.240^{+0.024}_{-0.025}$ | $0.289^{+0.031}_{-0.028}$ | $1.16^{+0.14}_{-0.11}$ | $0.205^{+0.016}_{-0.018}$ | $0.181^{+0.020}_{-0.019}$ |
| log g | Surface gravity (cgs) | 4.300 ^{+0.026} -0.036 | 3.987 ^{+0.030} -0.036 | 4.034 ^{+0.032} -0.033 | 4.374+0.039 | 3.940 ^{+0.022} -0.025 | 3.885 ^{+0.029} -0.031 |
| <i>T</i> _{eff} | Effective Temperature (K) | 6250 ⁺²²⁰ -190 | 6460 ⁺¹³⁰ -150 | 6274 ± 97 | -0.032 5597 ⁺⁸³ | -0.025 5948 ⁺⁸⁷ | -0.031 5910 ⁺⁹⁶ -99 |
| [Fe/H] | Metallicity (dex) | 0.258+0.078 | -150 $0.072^{+0.085}_{-0.076}$ | 0.110+0.078 | $0.078^{+0.072}_{-0.066}$ | $0.329^{+0.073}_{-0.075}$ | -99 0.299 ± 0.076 |
| $[Fe/H]_0^{\dagger}$ | | 0.258-0.080 0.271 ^{+0.070} -0.068 | 0.174 ^{+0.088} 0.174 ^{+0.088} | 0.212 ^{+0.083} 0.212 ^{-0.084} | 0.070 - 0.066 0.110 - 0.069 | 0.329 - 0.075 0.347 ± 0.069 | 0.296 ^{+0.070} 0.296 ^{-0.069} |
| | Initial Metallicity | | 0.174-0.080 | 0.212_0.084 2.33 ^{+0.71} 2.56 | | | 0.296-0.069 |
| Age | Age (Gyr) | 1.28+1.6 -0.91 | $1.99^{+0.55}_{-0.40}$ | .00 | 9.1 ^{+3.1} -3.9 | $2.64^{+0.38}_{-0.39}$ | $4.27^{+0.53}_{-0.42}$ |
| <i>EEP</i> [‡] | Equal Evolutionary Phase | 331+31 -37 | 380 ⁺¹⁹ -17 | 381+20 | 400^{+16}_{-37} | 403.1+7.9 | 452.9 ^{+4.3} -5.2 |
| $A_V \dots \dots$ | V-band extinction (mag) | $0.22^{+0.18}_{-0.14}$ | $0.217^{+0.055}_{-0.095}$ | _ | $0.062^{+0.039}_{-0.041}$ | $0.103^{+0.032}_{-0.053}$ | $0.086^{+0.043}_{-0.053}$ |
| σ_{SED} | SED photometry error scaling | $3.27^{+1.2}_{-0.73}$ | $0.84^{+0.35}_{-0.22}$ | _ | $0.70^{+0.29}_{-0.17}$ | $1.17^{+0.42}_{-0.27}$ | $1.20^{+0.44}_{-0.27}$ |
| σ | Parallax (mas) | 5.64 ± 0.10 | 2.925 ± 0.033 | _ | 6.536 ± 0.047 | 2.948 ^{+0.074} -0.075 | 2.973 ^{+0.077} -0.076 |
| d | Distance (pc) | $177.4^{+3.2}_{-3.1}$ | 341.8 ^{+3.9} -3.8 | _ | 153.0 ± 1.1 | 339.2 ^{+8.8} -8.3 | 336.4 ^{+8.9} -8.5 |
| lanetary Param | eters: | | -5.8 | | | -8.5 | -8.5 |
| P | Period (days) | $3.4095675^{+0.0000070}_{-0.0000069}$ | 5.0037775 ± 0.0000048 | 4.720219 ± 0.000011 | 10.180249 ± 0.000015 | 5.331751 ± 0.000011 | 5.331751 ± 0.0000 |
| R _P | Radius (<i>R</i> _I) | $1.060^{+0.041}_{-0.034}$ | $1.771^{+0.060}_{-0.056}$ | $1.396^{+0.056}_{-0.054}$ | $1.060^{+0.040}_{-0.039}$ | $1.239^{+0.046}_{-0.039}$ | $1.241^{+0.052}_{-0.049}$ |
| и Мр | Mass (<i>M</i> _I) | 6.33 ^{+0.29} -0.31 | -0.056 0.88 ± 0.16 | -0.054 2.37 ± 0.24 | $0.851^{+0.052}_{-0.047}$ | -0.039 0.99 ± 0.11 | -0.049 0.912 + 0.095 -0.10 |
| T_0^{\star} | Optimal conjunction Time (BJD _{TDB}) | $2458629.47972 \pm 0.00039$ | 2458450 72877+0.00071 | $2458913.37033 \pm 0.00045$ | 2458607.90338 ^{+0.00052} -0.00049 | 2458000 55202+0.00081 | 2458990.55306 ^{+0.00} -0.00 |
| 0 | | | $2458459.73877 + 0.00071 \\ -0.00083 \\ 0.06608 + 0.00098 \\ -0.0011 \\ 0.00$ | | 0.0903 ^{+0.0018} 0.0903 ^{+0.0018} | 2458990.55302 ^{+0.00081} 0.06864 ^{+0.00079} 0.06864 ^{+0.00079} | 2438990.33300 -0.00 |
| <i>a</i> | Semi-major axis (AU) | $0.04860^{+0.00080}_{-0.00094}$ | 0.06608-0.0011 | $0.0626^{+0.0011}_{-0.0012}$ | 0.0903-0.0013 | 0.06864-0.00075 | 0.06586+0.00069 |
| i | Inclination (Degrees) | 88.41 ^{+1.0} -0.93 | 82.54 ^{+0.42} -0.59 | 85.70 ^{+1.3} -0.65 | 88.51 ^{+0.29} -0.22 | 88.84 ^{+0.81} -1.1 | 88.4 ^{+1.1} -1.3 |
| e | Eccentricity | $0.072^{+0.021}_{-0.023}$ | 0.050 ^{+0.054} -0.035 | 0.073 ^{+0.092} -0.052 | $0.024^{+0.032}_{-0.017}$ | $0.036^{+0.044}_{-0.026}$ | $0.057^{+0.050}_{-0.039}$ |
| $\tau_{ m circ} \dots$ | Tidal circularization timescale (Gyr) | 3.32 ^{+0.53} -0.61 | $0.206^{+0.060}_{-0.057}$ | 1.29+0.37 -0.33 | $42.5^{+11}_{-8.1}$ | 1.85 ± 0.38 | $1.52^{+0.42}_{-0.38}$ |
| ω* | Argument of Periastron (Degrees) | -74^{+13}_{-15} | 159_99 | 105+64 | -110^{+120}_{-130} | -180^{+110}_{-100} | 120^{+57}_{-70} |
| T _{eq} | Equilibrium temperature (K) | 1586^{+52}_{-40} | $^{-99}_{1749^{+26}_{-30}}$ | 1679 ⁺³⁵ | -130 918 ± 11 | 1619^{+24}_{-23} | 1642^{+25}_{-24} |
| K | RV semi-amplitude (m/s) | 713+20 | -30 - 30 - 30 - 30 - 30 - 30 - 30 - 30 | 223+22 | 82.5 ^{+4.0} -3.8 | 87.3 ^{+9.2} | 87.6 ^{+8.9} 87.6 |
| $R_P/R_* \dots$ | Radius of planet in stellar radii | 0.08108+0.00075 -0.00071 | 0.08738 ^{+0.00091} -0.00086 | 0.0745 ^{+0.0014} -0.0015 | 0.1040 ± 0.0015 | 0.05827 ^{+0.00071} -0.00070 | 0.05834 ^{+0.00072} -0.00072 |
| - | - | -0.00071 | -0.00086 | -0.0015 c op+0.24 | | -0.00070 | |
| a/R* | Semi-major axis in stellar radii | 7.78+0.22 -0.27 | 6.82 ^{+0.22} -0.24 | 6.98 ^{+0.24} -0.23 | $18.54^{+0.70}_{-0.60}$ | 6.76 ^{+0.17} -0.20 | 6.48 ± 0.23 |
| Depth | Flux decrement at mid transit | 0.00657 ± 0.00012 | $0.00764^{+0.00016}_{-0.00015}$ | 0.00555 ± 0.00021 | 0.01083 ± 0.00032 | 0.003395+0.000083 -0.000081 | 0.003403+0.0000 |
| au | Ingress/egress transit duration (days) . | $0.01247^{+0.0012}_{-0.00061}$ | $0.0454^{+0.0028}_{-0.0026}$ | $0.0177^{+0.0027}_{-0.0033}$ | 0.0208 ± -0.0022 | $0.01483^{+0.00084}_{-0.00034}$ | 0.003403 -0.00008 0.01503 -0.00049 |
| <i>T</i> ₁₄ | Total transit duration (days) | $0.1576^{+0.0015}_{-0.0013}$ | 0.1502 ± 0.0017 | 0.1934+0.0025 | 0.1736 ± 0.0023 | 0.2627 ± 0.0020 | 0.2631+0.0022 |
| <i>b</i> | Transit Impact parameter | $0.23^{+0.13}_{-0.15}$ | $0.8763^{+0.0063}_{-0.0067}$ | 0.504+0.085 | $0.481^{+0.050}_{-0.076}$ | $0.136^{+0.13}_{-0.095}$ | 0.18+0.14 |
| $T_{S,14}$ | Total eclipse duration (days) | $0.1387^{+0.0063}_{-0.0055}$ | 0.8765_0.0067 0.1481 ^{+0.0070} 0.019 | 0.208+0.041 | 0.481 - 0.076 0.1736 + 0.0075 - 0.0066 | -0.095 0.263+0.014 0.013 | 0.270+0.024 |
| ρ _p | Density (cgs) | 6.58 ^{+0.70} 6.75 | 0.195 ^{+0.042} 0.040 | 1.08 ^{+0.18} -0.15 | 0.88 ^{+0.13} 0.88 ^{-0.11} | 0.640 ^{+0.097} 0.640 ^{-0.093} | -0.017 |
| | | 0.30-0.75 | 0.193-0.040 | 1.00-0.15 | 0.00-0.11 | 0.040-0.093 | $0.589_{-0.091}$ |
| $logg_P$ | Surface gravity | 4.145 ^{+0.031} -0.038 | 2.841 ^{+0.078} -0.095 | 3.479 ^{0.054} -0.056 | $3.273^{+0.045}_{-0.043}$ | 3.201 ^{+0.052} -0.059 | 3.165 ^{+0.055} -0.061 |
| $T_S \ldots \ldots$ | Time of eclipse (BJD _{TDB}) | 2458467.566 ^{+0.031} -0.038 | 2454824.46 ^{+0.12} -0.20 | 2458717.445 ^{+0.097} -0.13 | 2455701.431 ^{+0.092} -0.12 | 2454189.289 ^{+0.095} -0.17 | 2454189.27 ^{+0.12} -0.19 |
| $e\cos\omega_*$ | | $0.019^{+0.014}_{-0.018}$ | $-0.012^{+0.038}_{-0.062}$ | $-0.012^{+0.032}_{-0.042}$ | $-0.002^{+0.012}_{-0.018}$ | $-0.007^{+0.028}_{-0.049}$ | -0.013+0.035 -0.013+0.035 -0.057 |
| $e \sin \omega_* \ldots$ | | $-0.068^{+0.024}_{-0.022}$ | $0.007^{+0.042}_{-0.029}$ | $0.053^{+0.10}_{-0.056}$ | $0.000^{+0.030}_{-0.029}$ | 0.001+0.028 | 0.031+0.044 |
| | Separation at mid transit | -0.022 8.29 ^{+0.36} 8.29 ^{-0.40} | -0.029 6.76 ^{+0.38} 6.751 | 6.59 ^{+0.54} | -0.029 18.58 ^{+0.97} -0.88 | 6.74 ^{+0.32} -0.36 | $6.27^{+0.41}_{-0.47}$ |

NOTES:

See Table 3 in Eastman et al. (2019) for a detailed description of all derived and fitted parameters.

¶ No TESS dilution prior was used for TOI-1478 b because initial fits showed a fitted dilution past the 10% prior we used in the other systems. We fit for a dilution term within the fit for the *TESS* bandpass but with no prior.

[§] The SED was not included within the global fit for TOI-1333.

** T_C prior comes from analysis of the WASP photometry (see §2.2). We note that this time is in HJD_{TDB} while all data files and results here are BJD_{TDB}. The difference between these two time systems is on the order of seconds while the precision on T_C used as a prior is on order of minutes, and therefore has no influence on the results.

***Dilution prior for TOI-1333 comes from our 3-component SED analysis (see §2.6.4).

'We assume the TESS correction for blending is much better than 10%. We use a prior of 10% of the determined blending from TICv8 (Stassun et al. 2018). [†]The initial metallicity is the metallicity of the star when it was formed.

[‡]The Equal Evolutionary Point corresponds to static points in a stars evolutionary history when using the MIST isochrones and can be a proxy for age. See §2 in Dotter (2016) for a more detailed description of EEP.

*Optimal time of conjunction minimizes the covariance between T_C and Period. This is the transit mid-point.

^{π} The tidal quality factor (Q_S) is assumed to be 10⁶.

| Table 5. Median values and | l 68% confidence interva | ils for the global models |
|----------------------------|--------------------------|---------------------------|
|----------------------------|--------------------------|---------------------------|

| TOI-628 | | | | | | | | | |
|--|---|--|--|--|--|---|--|--|--|
| Wavelength 1 | Parameters: linear limb-darkening coeff | B 0.646 ^{+0.058} -0.061 | r' 0.329 ± 0.051 | z' 0.213 ^{+0.053} _0.052 | TESS 0.236 ^{+0.045} -0.046 | | | | |
| ^{<i>u</i>} ₁ | quadratic limb-darkening coeff | $0.040_{-0.061}$ 0.286 ± 0.057 | 0.329 ± 0.031 0.306 ± 0.049 | 0.213 ± 0.052 0.313 ± 0.050 | 0.292 ^{+0.046} 0.292 ^{+0.047} -0.048 | | | | |
| $u_2 A_D$ | Dilution from neighboring stars | - | - | - | $0.292_{-0.048}$ 0.0012 ± 0.0036 | | | | |
| Telescope Pa | | MINERVAF3 | MINERVAF4 | MINERVAF6 | TRES | | | | |
| γ_{rel} . | Relative RV Offset (m/s) | 19990 ⁺²⁶ -28 | 19969 ± 35 | 20032^{+26}_{-25} | -633^{+20}_{-18} | | | | |
| $\sigma_J \dots$ | RV Jitter (m/s) | 60 ⁺²⁹ -18 | 102^{+36}_{-24} | 83 ⁺²⁶ | 54 ⁺²⁸ 54 ²⁰ | | | | |
| σ_{J}^{2} | RV Jitter Variance | 3700 ⁺⁴⁴⁰⁰ -1800 | 10400^{+8600}_{-4300} | 7000^{+4900}_{-2900} | 2900 ⁺³⁸⁰⁰ -1700 | | | | |
| Transit Parar | neters: | TESS | KeplerCam UT 2019-12-06 (z') | LCO SSO UT 2019-12-16 (B) | Whitin UT 2020-02-08 (r') | | | | |
| $\sigma^2 \dots$ | Added Variance | $0.000000119^{+0.000000024}_{-0.000000020}$ | $0.0000366^{+0.0000020}_{-0.0000018}$ | $0.00001138^{\pm 0.0000010}_{-0.00000091}$ | $0.00000765^{+0.00000047}_{-0.00000044}$ | | | | |
| F_0 | Baseline flux | $0.999998^{+0.000039}_{-0.000038}$ | 0.99958 ± 0.00023 | 1.00071 ± 0.00021 | 1.00125 ± 0.00012 | | | | |
| <i>C</i> ₀ | Additive detrending coeff | — | _ | $0.00806^{+0.00058}_{-0.00057}$ | -0.00277 ± 0.00031 | | | | |
| TOI-640 | | | | 77 74 | | | | | |
| Wavelength | | I | z' | TESS | v | | | | |
| <i>u</i> ₁ | linear limb-darkening coeff | 0.191 ± 0.051 | 0.170 ± 0.036 | 0.216 ± 0.046 | 0.376 ± 0.050 | | | | |
| ^u 2 | quadratic limb-darkening coeff Dilution from neighboring stars | 0.311 ± 0.049 | 0.313 ± 0.035 | 0.314 ± 0.046 0.0000 ± 0.0022 | 0.329 ± 0.049 | | | | |
| A _D Telescope Pa | | CHIRON | _ | 0.0000 ± 0.0022 | - | | | | |
| | | | | | | | | | |
| γ_{rel} | Relative RV Offset (m/s) RV Jitter (m/s) | 39091^{+10}_{-11} $32.3^{+12}_{-8.3}$ | | | | | | | |
| $\sigma_J \dots \sigma_J^2 \dots$ | RV Jitter Variance | 1040^{+900}_{-470} | | | | | | | |
| Transit Parar | | 7040–470 TESS | Briefield UT 2019-12-27 (I) | PEST UT 2020-03-01 (V) | LCO SSO UT 2020-08-23 (z') | LCO SSO UT 2020-11-06 (z') | | | |
| σ ² | Added Variance | -0.000000009 ⁺ 0.000000093 -0.0000000081 | 0.00000259 ^{+0.00000059} -0.00000051 | 0.00000220 ^{+0.00000052} -0.0000045 | 0.00000187 ^{+0.00000053} -0.0000041 | 0.000000056 ^{+0.000000098} -0.000000080 | | | |
| F_0 | Baseline flux | -0.00000009 -0.000000081 1.000016 ± 0.000022 | 0.00000239 - 0.00000051 0.99995 ± 0.00018 | 0.00000220 - 0.00000045 0.99984 ± 0.00019 | 0.00000187 - 0.00000041 0.99956 ± 0.00023 | -0.00000000000000000000000000000000000 | | | |
| C_0 | Additive detrending coeff | | -0.0183 ± 0.0033 | - | 0.00080 ^{+0.00050} -0.00049 | -0.00052 ± 0.00025 | | | |
| C ₁ | Additive detrending coeff | _ | -0.0233 ± 0.0033 | - | -0.00049 | | | | |
| TOI-1333 | | | | | | | | | |
| Wavelength | Parameters: | I | R | z' | TESS | | | | |
| <i>u</i> ₁ | linear limb-darkening coeff | 0.228 ± 0.049 | 0.258 ± 0.046 | 0.183 ± 0.035 | $0.226^{+0.043}_{-0.044}$ | | | | |
| <i>u</i> ₂ | quadratic limb-darkening coeff | 0.323 ± 0.049 | 0.297 ± 0.048 | 0.303 ± 0.035 | 0.306 ± 0.047 | | | | |
| $A_D \dots$ | Dilution from neighboring stars | 0.345 ± 0.016 | 0.344 ± 0.017 | 0.0671 ± 0.0033 | -0.016 ± 0.028 | | | | |
| Telescope Pa | rameters: | TRES | | | | | | | |
| γ_{rel} | Relative RV Offset (m/s) | 122^{+15}_{-14} | | | | | | | |
| $\sigma_J \dots$ | RV Jitter (m/s) | 41 ⁺²¹ -18 | | | | | | | |
| $\sigma_J^2 \dots$ | RV Jitter Variance | 1700 ⁺²¹⁰⁰ -1200 | | | | | | | |
| Transit Parar | neters: | TESS | CRCAO UT 2020-07-29 (I) | LCO McD UT 2020-07-29 (z') | LCO McD UT 2020-08-12 (z') | CRCAO UT 2020-09-19 (R) | | | |
| $\sigma^2 \dots$ | Added Variance | $0.000000175^{+0.000000026}_{-0.000000023}$ | $0.00000078 \substack{+0.00000014 \\ -0.00000012}$ | $0.00000732^{+0.0000011}_{-0.00000097}$ | $\substack{0.00000353 + 0.00000052 \\ - 0.00000046}$ | $0.000000883^{+0.000000092}_{-0.000000084}$ | | | |
| F_0 | Baseline flux | 0.999984 ± 0.000036 | 0.99985 ± 0.00012 | 0.99927 ^{+0.00024} -0.00025 | 1.00245 ± 0.00018 | 1.00015 ± 0.00012 | | | |
| <i>C</i> ₀ | Additive detrending coeff | — | 0.00033 ± 0.00023 | -0.00060 ± 0.00065 | $0.01022\pm, 0.00068$ | -0.00017+0.00021 -0.00022 | | | |
| TOI-1478 | _ | _ | | | | | | | |
| Wavelength | | $R = 0.408 \pm 0.048$ | i' 0.349 ± 0.044 | TESS 0.410 ± 0.045 | | | | | |
| <i>u</i> ₁ <i>u</i> ₂ | linear limb-darkening coeff quadratic limb-darkening coeff | 0.408 ± 0.048 0.258 ± 0.049 | 0.349 ± 0.044 0.264 ± 0.047 | 0.320 ^{+0.048} | | | | | |
| A _D | Dilution from neighboring stars | - | - | $\begin{array}{c} 0.320 \substack{+0.048 \\ -0.047 \\ 0.126 \substack{+0.020 \\ -0.021 \end{array}}$ | | | | | |
| Telescope Pa | | CHIRON1 | CHIRON2 | -0.021 CORALIE | FEROS | TRES | | | |
| $\gamma_{\rm rel}$. | Relative RV Offset (m/s) | 19384.9 ± 7.0 | 19419 ± 11 | 20833.5 ^{+5.2} -5.4 | 20821.5 ± 4.5 | -30 ± 12 | | | |
| $\sigma_J \dots$ | RV Jitter (m/s) | 32.3 ^{+6.6} -5.5 | $1.5^{+37}_{-1.5}$ | $0.00^{+13}_{-0.00}$ | $10.4^{+4.2}_{-4.5}$ | $20.1^{+1.7}_{-3.3}$ | | | |
| $\sigma_J^2 \dots$ | RV Jitter Variance | 1040^{+470}_{-320} | 0^{+1500}_{-300} | -7^{+180}_{-91} | 108^{+110}_{-74} | 405 ⁺⁶⁹ -120 | | | |
| Transit Parar 2 | | TESS | KeplerCam UT 2019-12-14 (i') | PEST UT 2020-01-03 (R) | | | | | |
| σ^{2} | Added Variance | 0.000000083 ^{+0.000000049} -0.000000039 | 0.00000537 ⁺⁰ .00000039 -0.0000036 | 0.0000136+0.0000012 -0.0000011 | | | | | |
| F_0 | Baseline flux Additive detrending coeff | 1.000054 ± 0.000061 | 1.00003 ± 0.00016 0.00146 ± 0.00035 | 1.00026 ± 0.00022 | | | | | |
| TOI-1601 | . addute detrending coeff | | 0.00140 ± 0.00055 | | | | | | |
| Wavelength | Parameters: | I | R | z' | TESS | | | | |
| <i>u</i> ₁ | linear limb-darkening coeff | 0.281 ± 0.050 | 0.285 ± 0.050 | 0.215 ± 0.047 | 0.305 ± 0.045 | | | | |
| <i>u</i> ₂ | quadratic limb-darkening coeff | $0.305^{+0.050}_{-0.049}$ | $0.260^{+0.049}_{-0.050}$ | $0.281^{+0.049}_{-0.048}$ | 0.297 ± 0.048 | | | | |
| A _D | Dilution from neighboring stars | - | - | - | -0.0005 ± 0.0047 | | | | |
| Telescope Pa | | TRES | | | | | | | |
| γ_{rel} . | Relative RV Offset (m/s) | $-100.8^{+7.0}_{-7.3}$ | | | | | | | |
| $\sigma_J \dots = 2$ | RV Jitter (m/s) | $0.00^{+21}_{-0.00}$ | | | | | | | |
| $\sigma_J^2 \dots$ Transit Parar | RV Jitter Variance | -110+540 -300 TESS | GMU UT 2020-08-30 (R) | GMU UT 2020-09-15 (R) | CRCAO UT 2020-10-01 (R) | Adams UT 2020-10-17 (I) | LCO McD UT 2020-10-17 (z') | | |
| | Added Variance | -0.000000044 ^{+0.000000024} -0.000000020 | 0.00001324 ^{+0.0000010} -0.0000097 | 0.0000105 ^{+0.0000020} -0.000017 | 0.00000134 ^{+0.00000025} -0.0000021 | Adams UT 2020-10-17 (1) 0.00000066+0.00000017 -0.00000015 | | | |
| σ^2 | | -0.00000020 | | | -0.00000021 1.00043 ± 0.00012 | -0.00000000000000000000000000000000000 | $\begin{array}{c} 0.00000069 {+} 0.00000016 \\ {-} 0.00000014 \\ 0.998577 {+} 0.00091 \\ {-} 0.000092 \end{array}$ | | |
| $\sigma^2 F_0$ | Baseline flux | 0.999969 ± 0.000042 | 0.99318 ± 0.00018 | 1.00037 ± 0.00032 | 1.00043 ± 0.00012 | | 0.990.377 0 000000 | | |
| <i>F</i> ₀ | | | 0.99318 ± 0.00018 -0.01148 ± 0.00042 | -0.015 ± 0.00032 | -0.00039 ± 0.00012 | -0.00232 ± 0.00034 | -0.00039 ± 0.00024 | | |
| | Baseline flux | | | | | | -0.00039 ± 0.00024 -0.0015 ± 0.0013 | | |
| $F_0 \dots$ $C_0 \dots$ | Baseline flux | | -0.01148 ± 0.00042 | -0.015 ± 0.010 | | -0.00232 ± 0.00034 | -0.00039 ± 0.00024 | | |

sky-subtracted and flat-fielded. The reduced science frames were combined into a single combined image using an intrapixel interpolation that conserves flux, shifts the individual dithered frames by the appropriate fractional pixels, and median-coadds the frames (see Figure 4). The final resolution of the combined dither was determined from the fullwidth half-maximum of the point spread function; the resolutions of the $Br - \gamma$ and H - cont images are 0.092" and 0.075", respectively (Figure 4).

The sensitivities of the final combined AO images were determined by injecting simulated sources azimuthally around the primary target every 20° at separations of integer multiples of the central source's FWHM (Furlan et al. 2017, Lund et al. *in preparation*). The brightness of each injected source was scaled until standard aperture photometry detected it with 5σ significance. The resulting brightness of the injected sources relative to the target set the contrast limits at that injection location. The final 5σ limit at each separation was determined from the average of all of the determined limits at that separation and the uncertainty on the limit was set by the rms dispersion of the azimuthal slices at a given radial distance. The sensitivity curves for TOI-1333 are shown in Figure 4 along with an inset image zoomed to the primary target showing no other companion stars.

2.6.3. TOI-1333 Companions

In the case of TOI-1333, two additional sources were detected in the PHARO and GEMINI AO imaging (Figure 4, only the 3'' companion is shown). The first source is 7.43''(Gaia Collaboration et al. 2018) to the east of the primary target and is separately resolved by the 2MASS survey (2MASS J21400422+4824221; TIC 395171213). The second source is only 2.81" (Gaia Collaboration et al. 2018) away and was not separately detected by 2MASS - although it was detected by Gaia and hence is in the TESS Input Catalog (TIC 2010985858). Based upon the differential magnitudes measured by Palomar, the deblended infrared magnitudes for the primary star and nearby companion are $K_1 = 8.355 \pm 0.024$ mag, $H_1 = 8.514 \pm 0.043 \text{ mag} (H_1 - K_1 = 0.159 \pm 0.049)$ and $K_2 = 10.614 \pm 0.026$ mag, $H_2 = 10.871 \pm 0.044$ mag $(H_2 - K_2 = 0.257 \pm 0.051)$, respectively. The 3" star has a Gaia magnitude of $G = 12.6221 \pm 0.0016$.

The primary target has a Gaia distance of 200.5 ± 1.2 pc whereas the 7" companion has a Gaia distance of 1030 ± 35 pc indicating that this companion is not bound to the primary and is simply a chance alignment near the line-of-sight to the TOI-1333. However, the 3" companion has a Gaia distance of 196.8 ± 1.6 pc and proper motions that are nearly identical to that of the primary star ($\mu 1_{\alpha} = -9.81 \pm 0.05$ vs. $\mu 2_{\alpha} = -9.06 \pm 0.08$ and $\mu 1_{\delta} = -10.50 \pm 0.04$ vs. $\mu 2_{\delta} = -9.240.12$ mas/yr); it is, therefore, highly probable that TIC 2010985858 is physically bound to the primary star with a

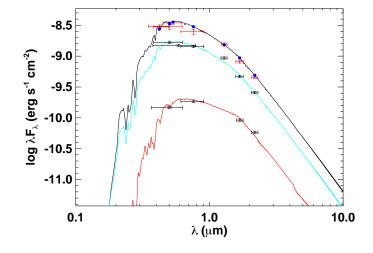


Figure 5. The three-component SED fit for TOI-1333. The blue points are the predicted integrated fluxes for the primary star. The red and black points are the observed values at the corresponding passbands for each star. The width of the passbands are the horizon-tal error bars and the vertical errors represent the 1σ uncertainties. The final model fit is shown by the solid line for TOI-1333 (black), and the 7" (cyan) and bound 3" (red) companions.

projected separation of \sim 470 au. We account for the blended light from these two companions in our global analysis (see §3 for details on how).

We use the LOFTI software package (Pearce et al. 2020) to derive orbital parameters of the visual binary system of the formed by TOI 1333 and its companion. LOFTI uses the relative proper motions of the two stars from the Gaia catalog to sample probable orbits for a binary star system. To derive the mass of the companion – required to fit orbit using LOFTI – we use the isochrones package (Morton 2015a). We perform an SED fit on photometry from TOI 1333's binary companion using the Gaia G, BP, and RP magnitudes, along with its Gaia parallax. Using the MIST isochrone (Dotter 2016) as the base isochrone, we derive a mass for the companion of $0.808^{+0.043}_{-0.042} M_{\odot}$. We use the astrometric parameters for the two systems from Gaia EDR3.

The LOFTI fit reveals that the semimajor axis of the binary orbit is 570^{+590}_{-170} AU and the orbital inclination is 125^{+18}_{-10} degrees, ruling out an edge-on orbit for the binary at high confidence. The orbital eccentricity is weakly constrained to be less than 0.69 with 95% confidence (with a slight preference for values between 0.5 and 0.7, but consistent with 0).

2.6.4. TOI-1333 Spectral Energy Distribution

The presence of the two stellar companions within a few arcseconds of TOI-1333 implies that the *TESS* and SG1 light curves of the TOI-1333 planet transit are likely to be diluted to some extent by the light from these other stars. Although the QLP corrects the *TESS* lightcurve for the blended contributions of known targets in the *TESS* input catalog (TIC),

0.05 0.04 Probability 0.03 0.02 0.01 0.00 5 1.3 1.4 1.5 2 3 4 7 1.2 1.6 1.7 1 6 Stellar Mass (M_o) Age (Gyr)

Figure 6. The (Left) M_{star} and (Right) Age probability distribution function for TOI-1601 from our global fit. We split this poster at the valley of $M_* = 1.415 \ M_{\odot}$, and extract two separate solutions, one for each of the peaks in the posteriors (see Table 4). The red line shows the median value for each parameter from the higher-mass solution with a probability of 68.4% (see Section 3).

we need to correct the follow-up photometry from TFOP for different amounts of dilution. To quantify this flux dilution, we performed a multi-component SED fit with Kurucz model atmospheres following the procedures described in Stassun & Torres (2016), utilizing the resolved broadband measurements from *Gaia*, 2MASS, and our AO observations (see §2.6.2).

We adopted the spectroscopic $T_{\rm eff}$ (6250±100 K) from TRES for TOI-1333 and the $T_{\rm eff}$ from the TICv8 (Stassun et al. 2019) and from the Gaia DR2 catalog for the companions, with A_V being left as a free parameter but limited to the maximum line-of-sight value from the dust maps of Schlegel et al. (1998). The resulting fits are shown in Figure 5, with best-fit $A_V = 0.06 \pm 0.06$ for the TOI-1333 planet host. Integrating the SED models over the TESS bandpass gives the total flux dilution $(F_2+F_3)/F_1 = 0.55$ with values 0.56 in the *I*-band 0.54 in the *R*-band, and 0.58 in the Sloan z'. We also used the SED fits to constrain the contribution from only the 2'' companion since some of our follow-up photometry only resolved the 7" companion. The flux dilution of F_2/F_1 is 0.07 in the I-band, 0.06 in the R-band, and 0.07 in the Sloan z'. By combining the deblended spectral energy distribution of the primary star TOI-1333 with the known Gaia DR2 parallax, we measure the its radius to be $R_{\star} = 1.963 \pm 0.064 R_{\odot}$. We use this as a prior on the global fit for TOI-1333 (see §3).

2.7. Location in the Galaxy, UVW Space Motion, and Galactic Population

For each of the TOIs analyzed here, we used their parallaxes, proper motions, and radial velocities and associated uncertainties from the Gaia DR2 catalog (Gaia Collaboration et al. 2018) to determine their location, kinematics, orbits, and associations with known stellar populations⁸. We corrected the native DR2 parallaxes and uncertainties using the prescription given in Lindegren et al. (2018). From these we computed the *prior* (DR2-based) estimates of the distances to the systems⁹. We used these to compute the heliocentric UVW space motions of the host stars, and then corrected for the sun's motion (UVW)_{\odot} with respect to the Local Standard of Rest (LSR) as determined by Coşkunoğlu et al. (2011). These resulting (UVW) values are shown in Table 1. We note that we adopt a right-handed coordinate system, such that positive *U* is toward the Galactic center.

We used the Galactic latitudes and distances to the systems to estimate their Z height relative to the sun, and then corrected for the $Z_{\odot} \simeq 30$ pc offset of the sun from the Galactic plane as determined by Bovy (2017) from the analysis of local giants. We use the UVW velocities relative to the LSR to determine the likelihood that the star belongs to thin disk, thick disk, halo, or Hercules stream¹⁰, using the categorization scheme of Bensby et al. (2014). We also report the estimates of the parameters of the Galactic orbits of the systems as determined by Mackereth & Bovy (2018) using Gaia DR2

⁸ We acknowledge that some of the analysis in this section was inspired by and follows that of Burt et al. (2020)

¹⁰ The Hercules stream is a chemically heterogeneous population of nearby stars that have distinct kinematics relative to the bulk of the local stellar distribution. See, e.g., Bensby et al. 2007.

⁹ We note that, for self-consistency, we explicitly did not adopt the *posterior* values of the parallaxes from the global fit as listed in Table 4.

astrometry and radial velocities¹¹. We estimated the spectral type of each TOI using its T_{eff} as determined from the global fit and given in Table 4, and using the T_{eff} -spectral type relations of Pecaut & Mamajek (2013). We then compared the position and orbits of the systems to the scale height h_Z of stars of similar spectral type as determined by Bovy (2017).

Finally, we also consider whether the systems may belong to any of the known nearby ($\lesssim 150$ pc young (\lesssim Gyr) associations using the BANYAN Σ (Bayesian Analysis for Nearby Young AssociatioNs Σ) estimator (Gagné et al. 2018). Not suprisingly, none of the systems had any significant ($\gtrsim 1\%$) probability of being associated with these young associations, and the BANYAN Σ estimator assigned all five systems as belonging to the 'field' with a high probability $\gtrsim 99\%$. We now discuss the results for each of the systems individually.

TOI-628: We find a distance from the sun of $d = 178 \pm 3$ pc, consistent with the *posterior* value listed in Table 4, and $Z - Z_{\odot} \simeq -23$ pc. We derive velocities relative to the LSR of $(U, V, W) = (-8.7 \pm 0.4, 3.2 \pm 0.2, 4.2 \pm 0.1)$ km s⁻¹. According to the categorisation scheme of Bensby et al. (2014), the system has a > 99% probability of belonging to the thin disk. The Galactic orbit as estimated by Mackereth & Bovy (2018) has a perigalacticon of $R_p = 7.67$ kpc, and apogalacticon of $R_a = 8.19$ kpc, an eccentricity of e = 0.03, and a maximum Z excursion from the Galactic plane of $Z_{\text{max}} = 63$ pc. This orbit is both consistent with the current location of the system, as well as the scale height of 97 pc for stars of similar spectral type (F7V). Indeed, TOI-628 is relatively dynamically 'cold,' for its spectral type. In other words, it has an orbit that is fairly close to that of the local LSR.

TOI-640: We find a distance from the sun of $d = 340 \pm 4$ pc, consistent with the *posterior* value listed in Table 4, and $Z - Z_{\odot} \simeq -76$ pc. We derive velocities relative to the LSR of $(U, V, W) = (-16.8 \pm 0.2, -16.7 \pm 0.4, -8.7 \pm 0.2)$ km s⁻¹. According to the categorisation scheme of Bensby et al. (2014), the system has a ~ 99% probability of belonging to the thin disk. The Galactic orbit as estimated by Mackereth & Bovy (2018) has a perigalacticon of $R_p = 6.28$ kpc, and apogalacticon of $R_a = 8.16$ kpc, an eccentricity of e = 0.13, and a maximum Z excursion from the Galactic plane of $Z_{\text{max}} = 150$ pc. This orbit is both consistent with the current location of the system, and suggest that the system is nearing its maximum excursion above the plane. It is also consistent as the scale height of 85 pc for stars of similar spectral type (F5.5V).

TOI-1333: We find a distance from the sun of $d = 200 \pm 2$ pc, and $Z - Z_{\odot} \simeq 19$ pc. We derive velocities relative to the LSR of $(U, V, W) = (23.0 \pm 0.1, -1.00 \pm 0.3, -6.0 \pm 0.1)$ km s⁻¹. According to the categorisation scheme of Bensby et al. (2014), the system has a ~ 99% probability of belonging to the thin disk. The Galactic orbit as estimated by Mackereth & Bovy (2018) has a perigalacticon of $R_p = 7.25$ kpc, and apogalacticon of $R_a = 8.32$ kpc, an eccentricity of e = 0.07, and a maximum Z excursion from the Galactic plane of $Z_{\text{max}} = 91$ pc. This orbit is consistent with the current location of the system. It is also consistent as the scale height of 97 pc for stars of similar spectral type (F7V).

TOI-1478: We find a distance from the sun of $d = 153 \pm$ 1 pc, and $Z - Z_{\odot} \simeq 67$ pc. We derive velocities relative to the LSR of $(U, V, W) = (-37.0 \pm 0.3, 26.4 \pm 0.4, 32.5 \pm 0.4, 32.5$ (0.2) km s⁻¹. According to the categorisation scheme of Bensby et al. (2014), the system has a $\sim 88\%$ probability of belonging to the thin disk, and a $\sim 12\%$ probability of belonging to the thick disk (and negligible probabilities of belonging to the halo or Hercules stream). The Galactic orbit as estimated by Mackereth & Bovy (2018) has a perigalacticon of $R_p = 7.71$ kpc, and apogalacticon of $R_a = 10.34$ kpc, an eccentricity of e = 0.14, and a maximum Z excursion from the Galactic plane of $Z_{\text{max}} = 650$ pc. Unfortunately, Bovy (2017) was unable to determine the scale height of stars of similar spectral type (G6V) due to incompleteness. Nevertheless, it would appear that TOI-1478's orbit has a maximum Z excursion that exceeds the expected scale height for stars of similar spectral type as estimated by extrapolating from the results of Bovy (2017) from earlier spectral types. Surprisingly, its current distance above the plane is only small fraction of its predicted maximum excursion. In summary, the weight of evidence suggests that TOI-1478 may well be a thick disk star that we happen to be observing when it is near the Galactic plane. Detailed chemical abundance measurements (e.g., $\left[\alpha/\text{Fe}\right]$) may provide corroborating evidence for or against this hypothesis.

TOI-1601: We find a distance from the sun of $d = 336 \pm$ 9 pc, consistent with the *posterior* value listed in Table 4, and $Z-Z_{\odot} \simeq -73$ pc. We derive velocities relative to the LSR of $(U,V,W) = (-8.1\pm0.7, -14.5\pm0.7, 20.9\pm0.4)$ km s⁻¹. According to the categorisation scheme of Bensby et al. (2014), the system has a ~ 98% probability of belonging to the thin disk. The Galactic orbit as estimated by Mackereth & Bovy (2018) has a perigalacticon of $R_p = 6.50$ kpc, and apogalacticon of $R_a = 8.32$ kpc, an eccentricity of e = 0.12, and a maximum Z excursion from the Galactic plane of $Z_{\text{max}} = 351$ pc. This orbit is consistent with the current location of the system. The maximum Z excursion is a factor of ~ 3.3 times larger than the scale height of 103 pc for stars of similar spectral type (GOV). The probability that a star in a population with a given scale height $h_z = 108$ pc has a maximum

¹¹ We note that Mackereth & Bovy (2018) adopted a solar Galactocentric distance of $R_0 = 8$ kpc, $Z_{\odot} = 25$ pc, and a local circular velocity of $V_{\rm circ} = 220$ km s⁻¹. The also corrected for the sun's motion with respect to the LSR using the values of (UVW)_☉ as determined by Schönrich et al. (2010), which differ slightly from the values we adopted as determined by Coşkunoğlu et al. (2011). However, this is a minor effect and has no qualitative impact on our conclusions.

excursion of $z_{\text{max}} = 351$ pc is non-negligible. Thus we expect that TOI-1601 is a thin disk star that is simply in the tail of the distribution of z_{max} . Again, detailed abundances could corroborate or refute this conclusion.

3. EXOFASTv2 GLOBAL FITS

We use the publicly available exoplanet fitting suite, EXOFASTv2 (Eastman et al. 2013, 2019), to globally fit the available photometry and RVs to determine the host star and planetary parameters for TOI-628 b, TOI-640 b, TOI-1333 b, TOI-1478 b, and TOI-1601 b. We fit the TESS and SG1 transits (see $\S2.4$), accounting for the 30 min smearing from the FFIs. Within the fit, the SG1 lightcurves were detrended (additive) against the corresponding parameters shown in Table 1. See §D in the appendix of Collins et al. (2017) for a description of each detrending parameter. We use the MESA Isochrones and Stellar Tracks (MIST) stellar evolution models (Paxton et al. 2011, 2013, 2015; Choi et al. 2016; Dotter 2016) and the spectral energy distribution (SED) within the fit to determine the host star parameters for all systems but TOI-1333 b. The SED fit within the global fit puts systematic floors on the broadband photometry errors (see Table 1, Stassun & Torres 2016). We also note that EXOFASTv2 defaults a lower limit on the systematic error on the bolometric flux $(F_{bol} \sim 3\%)$ given the spread seen from various techniques to calculate it (Zinn et al. 2019). We place a Gaussian prior on the metallicity from our analysis of the host star's spectra from TRES, or CHIRON in the case of TOI-640 (see Section 2.5.3 and 2.5.2). We also place a Gaussian prior on the parallax from Gaia (Gaia Collaboration et al. 2016, 2018), correcting for the 30 μ as offset reported by Lindegren et al. (2018), and an upper limit on the line of sight extinction from Schlegel et al. (1998) & Schlafly & Finkbeiner (2011). We also fit for a dilution term on the TESS band. Since the QLP corrects the TESS light curves for all known blended companions, we place a Gaussian prior of $0\pm10\%$ of the contamination ratio reported by the TESS Input Catalog (TIC, Stassun et al. 2018). We assume that the light curve has been corrected to a precision better than 10% (and test this with preliminary EXOFASTv2 showing the dilution to be consistent with zero), but this flexibility also provides an independent check on the correction applied and allows us to propagate the uncertainty in the correction. We do not find any significant additional dilution in any of the systems (within this prior and consistent with zero dilution) other than TOI-1478 b, where our fit suggests an additional 13% dilution $(0.126^{+0.020}_{-0.021})$ is needed for the *TESS* light curve to be consistent with the TFOP photometry. For this fit, we remove this prior, essentially allowing the TFOP observations to constrain the depth of the transit. It is not clear what the cause of this additional dilution is since we see no evidence for any unknown companions in our high-spatial resolution imaging.

We note that TESS only observed three transits in one sector for TOI-1478 b, the longest period planet in our sample, and the TFOP light curves were both at higher spatial resolution. We also ran a fit where we allowed for a slope in the RVs, but found no significant trends for any system (we do not fit for a slope in the final fits). A list of the priors for each target is shown in Table 4. Table 3 of Eastman et al. (2019) shows a list and description of the fitted and derived parameters, including the bounds that EXOFASTv2 adopts for each fitted parameter. We note that eccentricity, a key parameter for this study, is bound as such: $0 \le e \le 1 - \frac{a+R_p}{R_*}$, in order to ensure that the periastron values of the planet orbits are larger than the sum of the stellar and planetary radius. We deem a fit to be fully converged by following the recommending statistical threshold of a Gelman-Rubin statistic (<1.01) and independent draw (>1000) that is recommended by Eastman et al. (2019). The results from our EXOFASTv2 fits are shown in Table 4 & 5, and the models are shown for the transits and RVs in Figures 2 & 3.

In the case of TOI-1333 b, we deviate slightly from the methodology in the previous paragraph because there are two nearby bright companions, both detected by high resolution imaging (see §4). The 2'' nearby companion and 7'' star were blended in the *TESS* and CRCAO photometry (see $\S2.4$), but only the 2" companion was blended in the LCO observations. While the TESS light curve has already been deblended as part of the reduction pipeline (see §2.1), the SG1 observations were not. Our three-component SED analysis (see §2.6.4) determined that the nearby companion 2'' from TOI-1333 accounts for 6.6% in the z'-band, where LCO did not resolve the close companion. The combined flux contribution from both stars is 36.1% in the I-band and 35.0% in the R-band, where both companions were unresolved by CRCAO. We use these values with a 5% Gaussian prior EXOFASTv2 global fit (also placing the prior on the TESS dilution as discussed in the previous paragraph). We place a Gaussian prior on the host star's radius from the SED analysis of $R_{\star} = 1.963 \pm 0.064 R_{\odot}$. Preliminary SED fits of TOI-1333 using EXOFASTv2 and independent log g constraints from the SPC analysis of the TRES spectra suggested that TOI-1333 is a slightly evolved star. Given that the SED would normally constrain the $T_{\rm eff}$ of the host star within the fit but was excluded for TOI-1333, we also place a prior on the $T_{\rm eff}$ of 6250±100 K from the SPC analysis of the TRES spectroscopy.

For TOI-1478 global fit, five RVs were acquired in the summer and winter of 2020, after a multi-month shutdown due to the COVID-19 pandemic. When included as part of the CHIRON RVs in the fit, we see a statistically significant slope measured of 0.176 m/s/day. However, the RV baseline for CHIRON observations has appeared to shift slightly when pre- and post-shutdown RVs were compared for stan-

dard stars, consistent with the shift we measured when including these post pandemic RVs (T. Henry, private communication). Since we saw no evidence of a slope in fits without these RVs and we know there is a shift observed in the RV baseline for CHIRON, we treat these observations as a separate instrument (labeled "CHIRON2"). When fitting the CHIRON RVs separately within the EXOFASTv2 fit, we see a difference of 65 m s⁻¹ in the fitted RV zero points. We note that the planet and host star parameters are consistent to <1 sigma whether or not this slope is included in the fit.

3.1. TOI-1601 Bimodality

After each EXOFASTv2 fit, we inspect the posteriors of each fitted and derived parameter, visually inspecting for any anomalies such as multi-modal distributions. In all cases but TOI-1601 b, no issues were noted. For TOI-1601 b, we see a clear bimodal distribution in the mass and age of the host star (see Figure 6). We find two peaks in the mass distribution at 1.340 and 1.517 M_{\odot} , which corresponds to the two peaks seen in the age distribution at 4.27 and 2.63 Gyr. There is no optimal way to represent the bimodal solution, so we split the mass of the host star at the minimum value of $1.415 M_{\odot}$ and extract two solutions for both peaks identified. We adopt the high mass solution for the discussion since it has a higher probability (66.7%) of being correct from our analysis, but present both solutions in Table 4 for future analysis. We note that we observed no significant change in any systematic parameters, suggesting the bimodality to be due to our limited precision that is not sufficient to completely separate similar solutions due to the host star being slightly evolved.

4. DISCUSSION

These five newly-discovered hot Jupiters from NASA's TESS mission significantly increase (>5%) the number of well characterized (measured masses and orbital eccentricities) giant planets that reside in orbits with periods between 5 and 15 days, a regime where planets experience weaker tidal forces than those experienced by planets closer to their host stars, and as a result, are not likely to have had enough time to circularize their orbits. Additionally, these longer period systems enable us to explore the "reinflation" scenario, an area TESS should make a huge impact in given its high photometric precision. By continuing to discover and characterize new giant planets at longer periods, we can look for evidence (through their eccentricity distribution) of the dominant migration mechanism. Additionally, TESS will provide a complete magnitude-limited sample of hot Jupiters (P<10 days, Zhou et al. 2019), allowing us to test whether multiple populations exist within the distribution of key parameters (mass, radius, eccentricity), where some tentative trends have been suggested (Nelson et al. 2017: Ikwut-Ukwa et al. 2021). Here we provide a short overview of our global fit results on each of the five new systems. We note that all three

of the planets found to be orbiting subgiant host stars (TOI-640, TOI-1333, and TOI-1601) reside in circular orbits.

4.1. TOI-628 b

TOI-628 (TIC 281408474) is a V=10.18 late-F star with a mass of $M_* = 1.311_{-0.075}^{+0.066} M_{\odot}$, radius of $R_* = 1.345_{-0.040}^{+0.046} R_{\odot}$, and an age of $1.28_{-0.91}^{+1.6}$ Gyr. Its planetary companion (TOI-628 b) has a radius of $R_{\rm P} = 1.060^{+0.041}_{-0.034} R_{\rm J}$, a mass of $M_{\rm P} = 6.33^{+0.29}_{-0.31} M_{\rm J}$, and is on a 3.4096-day period orbit. Our global analysis measures a non-zero orbital eccentricity of e = $0.072_{-0.023}^{+0.021}$. Within our global fit, we derive a circularization timescale of $\tau_{\rm circ} = 3.32^{+0.53}_{-0.61}$ Gyr (for this system, assuming equilibrium tides and tidal quality factors for the planet and star of Q_p 10⁶ and Q_* 10⁶ Adams & Laughlin 2006) for this system, which is longer than our estimated age from MIST of $1.28^{+1.6}_{-0.91}$ Gyr. Thus the small but non-zero eccentricity is likely a vestige of the initially high eccentricity that the planet obtained during some process that initiated high-eccentricity migration, a high eccentricity that was subsequently damped to the eccentricity we see today. Without a tighter constraint on the system's age, this is not conclusive. Also of interest is the high mass of TOI-628 b, which makes it one of only a few dozen known hot Jupiters with a mass $>6 M_{\rm J}$, and the most massive hot Jupiter found from TESS to date (see Figure 7).

4.2. TOI-640 b

The host star TOI-640 (TIC 147977348) is an F-star with a mass of $1.526^{+0.072}_{-0.079} M_{\odot}$ and a radius of $2.082^{+0.064}_{-0.058} R_{\odot}$. The host star appears to be just transitioning off the main sequence into the subgiant branch as suggested by our measured $\log g = 3.987^{+0.030}_{-0.036}$ dex (cm s⁻²) and corresponding tight age constraint within our global fit from the MIST evolutionary tracks of $1.99^{+0.55}_{-0.40}$ Gyr. It hosts a planetary companion, TOI-640 b, which is highly inflated ($R_{\rm P} = 1.771^{+0.060}_{-0.056}$ $R_{\rm J}$) Jupiter mass ($M_{\rm P} = 0.88 \pm 0.16 M_{\rm J}$) planet with a near integer orbital period of 5.0037775 ± 0.0000048 days. The orbit of the planet is consistent with circular, $e = 0.050^{+0.054}_{-0.035}$. It is only the third hot Jupiter known with a highly inflated radius (R_P >1.7) and on a period >5 days, joining KELT-12 b (Stevens et al. 2017) and Kepler-435 b (Almenara et al. 2015). Interestingly, TOI-640 b is almost a twin of KELT-12 b, in that they are highly inflated Jupiter-mass planets on \sim 5 day orbits around similar subgiant host stars. All three host stars in this regime are evolved, possibly suggesting that the inflation is a result of the host star's recent evolution (Assef et al. 2009; Spiegel & Madhusudhan 2012; Hartman & Bakos 2016; Lopez & Fortney 2016). Similarly to KELT-12b and Kepler-435 b, we see no evidence of any significant eccentricity.

TESS Mp > 0.4 e = 0 NEA 0.6 2.0 6500 TESS $\sigma e > 0$ 0.5 1.8 6000 Orbital Eccentricity Planet Radius [R_l] 1.6 ²²⁰⁰ T_{eff} (K) M_P (M_I) $\overline{\mathbf{\Phi}}$ 1.4 5000 4500 0.1 1.0 4000 0.0 10 10 Log(Orbital Period [Days]) Log(Orbital Period [Days])

Figure 7. (Left) The eccentricity and log of the orbital period of all known giant planets with a mass greater than $0.4M_J$ with period between 0.8 and 16 days. The *TESS* discovered systems are colored by the host star's effective temperature. The systems with a measured eccentricity from the NASA Exoplanet Archive (NEA) are shown as black circles with errors. Systems where the eccentricity was assumed to be zero are shown with gray crosses. (Right) The radius and log of the orbital period of all known transiting giant planets. The systems known prior to *TESS* are in black, while the systems discovered by TESS, including those presented in this paper and Ikwut-Ukwa et al. (2021), are shown as circles colored by their planet's mass.

TOI-1333 (TIC 395171208) is a bright (V=9.49) evolved F-star with a mass of $M_* = 1.464^{+0.076}_{-0.079} M_{\odot}$ and radius of $R_* = 1.925^{+0.064}_{-0.063} R_{\odot}$. The star appears to be slightly evolved, as suggested by its log g of $4.034^{+0.032}_{-0.033}$ dex (cm s⁻²). As a result of its evolutionary stage, we estimate a relatively tight age constraint from the MIST evolutionary tracks of $2.33^{+0.71}_{-0.56}$ Gyr. Orbiting on a 4.720219 ± 0.000011 day period, the planetary companion TOI-1333 b has a radius of $1.396^{+0.054}_{-0.054} R_J$, a mass of $M_P = 2.37 \pm 0.24 M_J$, and an eccentricity that is consistent with circular (e = $0.073^{+0.092}_{-0.052}$). This is not surprising given that our derived circularization timescale, $1.29^{+0.37}_{-0.33}$ Gyr, is similar to the age of the system.

In the case of TOI-1333, we have measured a periodicity of \sim 5.3 days from the ground-based and TESS photometry (we note that WASP identified a period $3 \times$ this). We have also measured a $v \sin I_*$ of 16.5±0.5 km s⁻¹. If the periodicity identified in the photometry indeed is the average rotation period of the host star, then we can estimate the inclination of the host star's rotation axis and compare it to the derived inclination of TOI-1333 b's orbit following the methodology presented in Masuda & Winn (2020). Using the EXOFASTv2 implementation of a Markov Chain Monte Carlo, we run a simple fit of the host star's rotational velocity and its projection onto our line of sight $(v \sin I_*)$ using the values from our global fit for Rstar $(1.925^{+0.064}_{-0.063})$ R_{\odot}), the derived rotational period of TOI-1333 from TESS and KELT (5.3 days), and the $v \sin I_*$ from the TRES spectroscopy $(14.2\pm0.5 \text{ km s}^{-1})$ to calculate the inclination of TOI-1333's rotation axis (relative to our line of sight). The latitudes on the Sun that show star spots have a differen-

tial rotation on the surface of a few percent. Therefore, we place a 3% error on the rotational velocity for this analysis. We require the same Gelman-Rubin statistic (<1.01) and independent draw (>1000) for convergence as the default for EXOFASTv2. We derive the inclination of the rotation axis to be $51.3^{\circ+3.5^{\circ}}_{-3.3^{\circ}}$. From our global fit, TOI-1333 b has an inclination of $85.70^{\circ+1.3^{\circ}}_{-0.65^{\circ}}$, suggesting that the rotation axis of the star and the orbital plane are misaligned. TOI-1333 b is an excellent candidate to confirm this result through spinorbit alignment (λ) measurements using the R-M or Doppler tomography techniques. The planet's orbit is also misaligned with the orbit of the wide binary companion TOI-1333 B, for which we measured an inclination of 125^{+18}_{-10} degrees from our LOFTI analysis. Interestingly, we do not detect a significant orbital eccentricity from our global fit for TOI-1333 b (though a small eccentricity is still possible) but this suggested misalignment might be a remnant left over from high eccentricity migration. The likely bound companion at 470 au (see §2.6.2) could be responsible for Kozai-Lidov migration of the planet.

4.4. TOI-1478 b

TOI-1478 (TIC 409794137, V = 10.81) is a Sun-like Gdwarf with radius of $R_* = 1.048^{+0.030}_{-0.029} R_{\odot}$, mass of $M_* = 0.946^{+0.059}_{-0.041} M_{\odot}$, and an age of $9.2^{+3.1}_{-3.9}$ Gyr. Orbiting TOI-1478 is a warm-Jupiter with a period of 10.180249 ± 0.000015 days, a radius of $R_{\rm P} = 1.060^{+0.040}_{-0.039} R_{\rm J}$, a mass of $M_{\rm P} = 0.851^{+0.052}_{-0.047} M_{\rm J}$, and resides in a circular orbit (e = $0.024^{+0.032}_{-0.017}$). TOI-1478 b is the longest period planet in our sample, and the planet and its host star (other than their orbital distances) resemble the Sun and Jupiter in mass and radius, possibly an example of an alternate outcome of our own Solar System. As a result of the long orbital period, the tidal forces on TOI-1478 b are too weak to have circularized the orbit. Therefore, the lack of a significant eccentricity could suggest a more dynamically quiescent migration history.

4.5. TOI-1601 b

In the case of TOI-1601 b, our global model showed a clear bimodality in the posterior distribution of the host star's mass and age (See $\S3.1$). This is likely due to the host star's evolutionary status, because the star sits on the HR diagram where isochrones cross, so the evolutionary state is ambiguous given the precision of our observations. and to account for this we extract two separate solutions, one for each peak in our posteriors. The higher host star mass solution, $M_* =$ $1.517_{-0.049}^{+0.053} M_{\odot}$, has a higher probability of being correct at 66.7%, so we adopt this solution for the discussion but both results are available in Table 4. TOI-1601 (TIC 139375960. V = 10.71) is an evolved subgiant (log g of $3.940^{+0.022}_{-0.025}$ dex V = 10.71) is an evolved subgrant (1058 of 20.76 $_{-0.025}$ (cm s⁻²)) with a radius of $R_* = 2.186^{+0.074}_{-0.063} R_{\odot}$. We estimate the age of the system within our fit to be $2.64^{+0.38}_{-0.39}$ Gyr. TOI-1601 b is a Jupiter mass planet $(0.99 \pm 0.11 M_J)$ that shows some inflation $(R_{\rm P} = 1.159^{+0.062}_{-0.059} R_{\rm J})$ and a circular orbit (e $= 0.037^{+0.045}_{-0.026}$), 5.331752 ± 0.000011 day orbit. The spectroscopic analysis of the TRES spectra of TOI-1601 shows some metal enhancement ([Fe/H] = $0.316^{+0.072}_{-0.074}$).

4.6. TESS's impact on Giant Planets

While the primary goal of NASA's TESS mission is to discover and measure the masses of small planets (Ricker et al. 2015), TESS has already provided some valuable discoveries in the field of giant planets (see, e.g., Huang et al. 2020b; Armstrong et al. 2020; Vanderburg et al. 2020). Given the minimum \sim 27 day baseline for any target, and the completeness in the sensitivity of space-based photometry to detect a hot-Jupiter transit, TESS provides the opportunity to obtain a near-complete sample of hot Jupiters (Zhou et al. 2019). To date, TESS has discovered 26 giant planets ($M_{\rm P} > 0.4 M_{\rm J}$), 16 of which have an orbital period >5.0 days (these numbers include the five systems presented in this paper and 2 additional systems from (Ikwut-Ukwa et al. 2021). For comparison, 36 hot Jupiters have been discovered with orbital periods >5.0 days from ground-based transit surveys (NASA Exoplanet Archive, Akeson et al. 2013).

If giant planets predominantly migrate through dynamical interactions, we may find evidence of this evolutionary history in the eccentricity distribution of hot Jupiters, specifically those that are dynamically young (where the circularization timescale by tidal forces is longer than the age of the system). Figure 7 shows the current distribution of giant planet eccentricities as a function of orbital period out to

16 days. Those 26 systems discovered by TESS are shown colored by their host star's effective temperature. Although this is not a homogeneous sample, since a variety of different analysis methods and assumptions were made within this population, there is a wider distribution of eccentricities for those systems with an orbital period > 5 days, where tidal circularization timescales are longer (Adams & Laughlin 2006). Interestingly, of the 5 systems presented here, only TOI-628 b has a statistically significantly measured eccentricity (e = $0.072^{+0.021}_{-0.023}$), and is consistent with dynamically driven migration since its estimated age is less than the circularization time-scale of the orbit. Although the other systems show some non-zero eccentricities from our global fits, they are not statistically significant (>3 σ) and could be a result of the Lucy-Sweeney bias (Lucy & Sweeney 1971). We also note that there is one very massive hot Jupiter in our sample, TOI-628 b (M_P 6.33 $^{+0.29}_{-0.31}$ M_J), and it is the most massive hot Jupiter discovered to date by TESS (we note that TESS has discovered a few transiting brown dwarfs, Jackman et al. 2019; Šubjak et al. 2020; Carmichael et al. 2020b,a and WD 1856+534 which has a mass limit $<13.8 M_{\rm J}$, Vanderburg et al. 2020). These massive Jupiters provide a great laboratory to study the effect of high gravity on the atmosphere of a gas giant, while studying the transition point between giant planets and brown dwarfs.

5. CONCLUSION

We present the discovery and characterization of five new giant planets (TOI-628, TOI-640 b, TOI-1333 b, TOI-1478 b, and TOI-1601 b) from NASA's TESS mission. These planets were discovered in the primary mission using the 30-minute cadenced, full frame images. Of the systems TOI-640 b, TOI-1333 b, and TOI-1601 b all orbit stars that appear to have just evolved off the main sequence entering the subgiant phase, as suggested by their estimated $\log g$ being under 4.1 dex (cm s^{-2}). None of the planets orbiting these subgiants appear to reside in a significantly eccentric orbits. TOI-628 b is the most massive hot Jupiter discovered by TESS ($M_{\rm P}$ = $6.33^{+0.29}_{-0.31}$ M_J), and resides in an eccentric orbit that is consistent with dynamically driven migration. Another planet from this work, TOI-640 b, is one of the only highly inflated $(> 1.7 R_{\rm J})$ hot Jupiters with an orbital period greater than 5 days. TOI-1478 b is the only planet in this sample with an orbital period >10 days, and it and its star are similar in size and mass to Jupiter and the Sun. All five planets orbit bright (V < 10.7) host stars and significantly increase the sample of well-characterized, long period (>5 day) hot Jupiters, an area where NASA's TESS mission should continue to provide a wealth of discoveries.

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